

NIMROD simulations of flux injection in a coplanar gun

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Inductive Flux Injection Model

- NIMROD has two available flux injection models^a
 - direct flux injection via Faraday’s law (specifying tangential \mathbf{E} field at the boundary)
 - inductive flux injection via Ampere’s law (specifying a tangential \mathbf{B} field at the boundary)
- inductive flux injection (i.e. specifying input current) is closer to (coaxial gun source) experiment
- exploit integral Ampere’s law $\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 I$
- specify B_ϕ at the boundary corresponding to coaxial gap to induce poloidal current in simulation domain

^aC. Sovinec, “Ohmic Current Drive in NIMROD Simulation”, NIMROD internal note, 2005

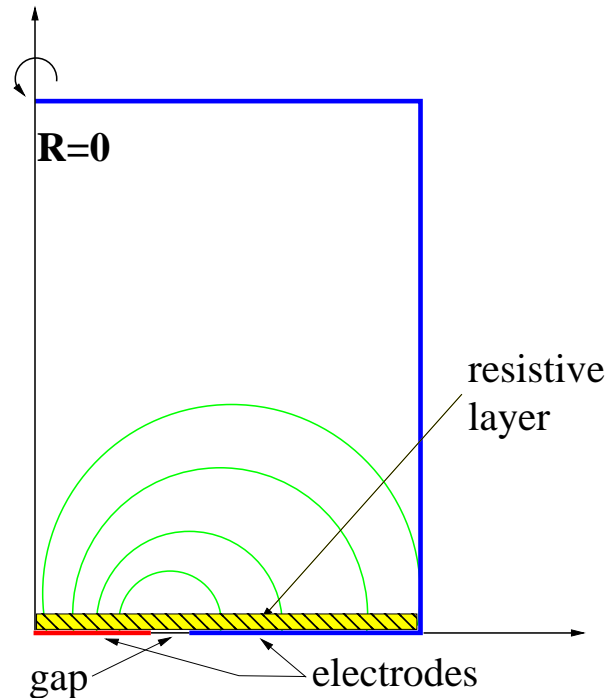


Inductive Flux Injection Model cont.

- along the boundary corresponding to flux gap specify $B_\phi R$
 - otherwise boundary condition is no-slip perfect conductor
- prescribed $B_\phi R$ induces a poloidal current in the simulation domain
- amplitude of $B_\phi R$ may vary with time (e.g. constant slope or programmed from experiment)
- ? thin highly resistive layer along bottom to rapidly diffuse flux across bottom
- the resulting $\mathbf{J} \times \mathbf{B}$ force pulls in the flux, drives columnation

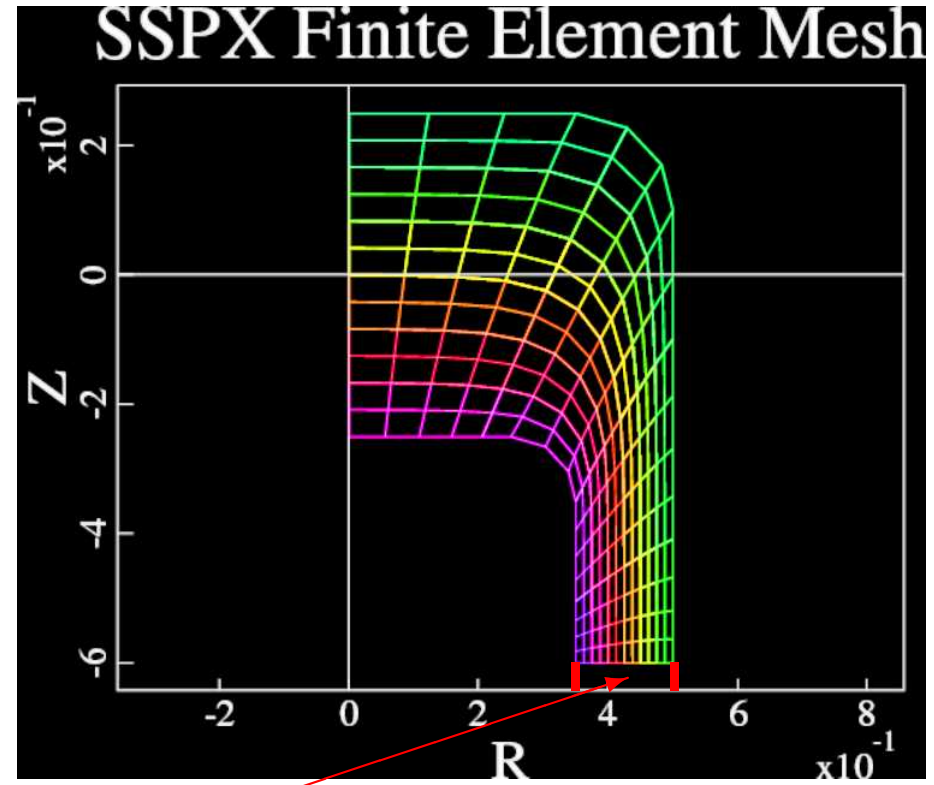
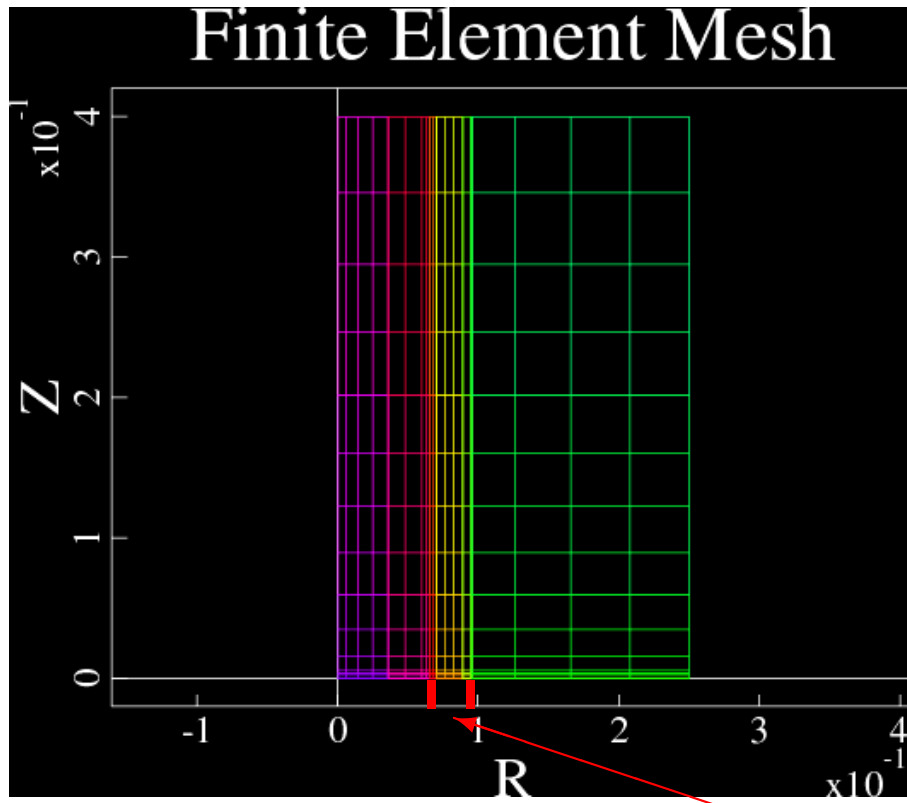
Coplanar Flux Injection Simulation

- cylindrical vessel with small flux gap of a few centimeters



- vacuum field is dipole-like $\sim .1mWb$
 - flux gap located at top of the arc
- peak current is $40 - 100kA$ ramped over $4 - 15\mu s$
- thin highly resistive layer across bottom 10^5 larger than background resistivity

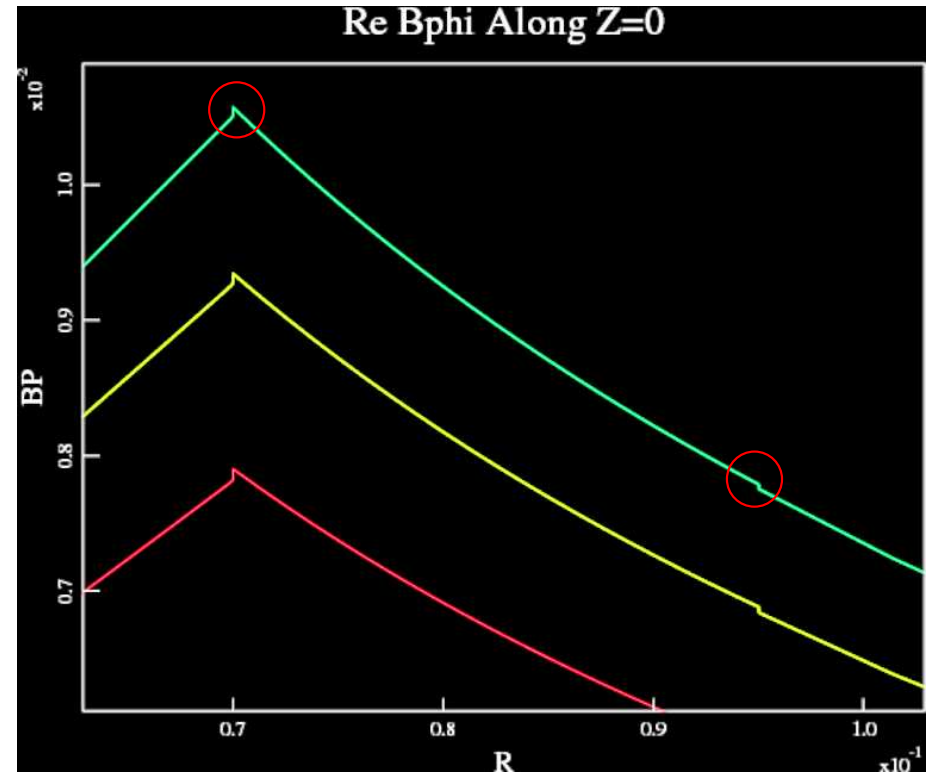
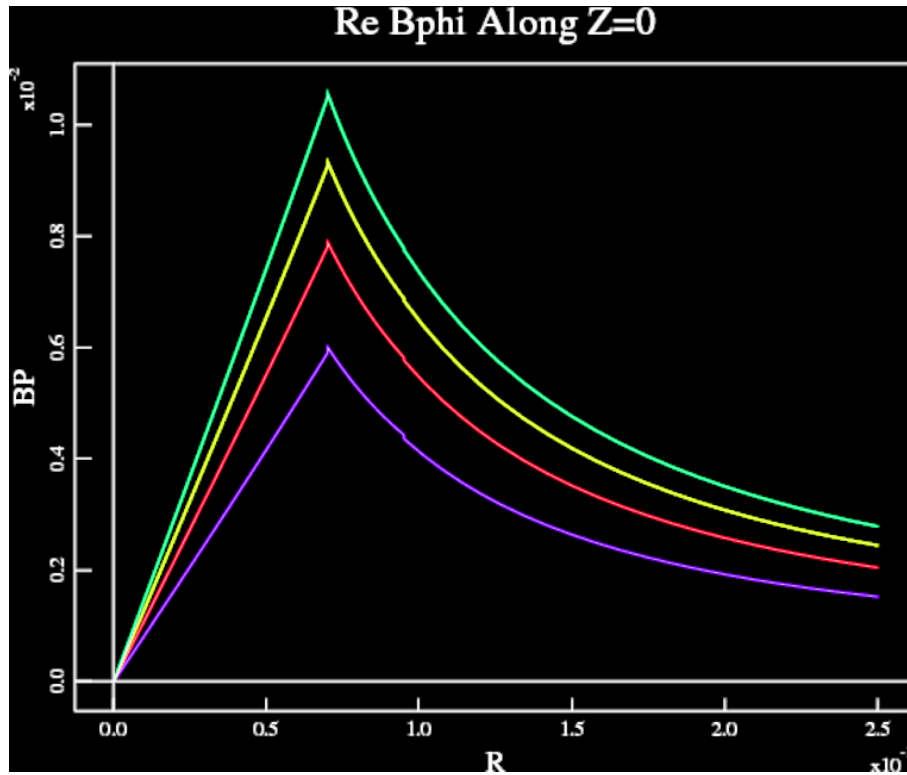
Coplanar Flux Injection is more challenging than coaxial injection



flux gap

- larger inductive \mathbf{E} field
- coaxial gun is well resolved
- discontinuous jump in $\frac{\partial B_\phi}{\partial R}$
- $\lambda_g^{SSPX} \sim 10 - 15m^{-1}$ - we attempt $\lambda_g \sim 10 - 100m^{-1}$

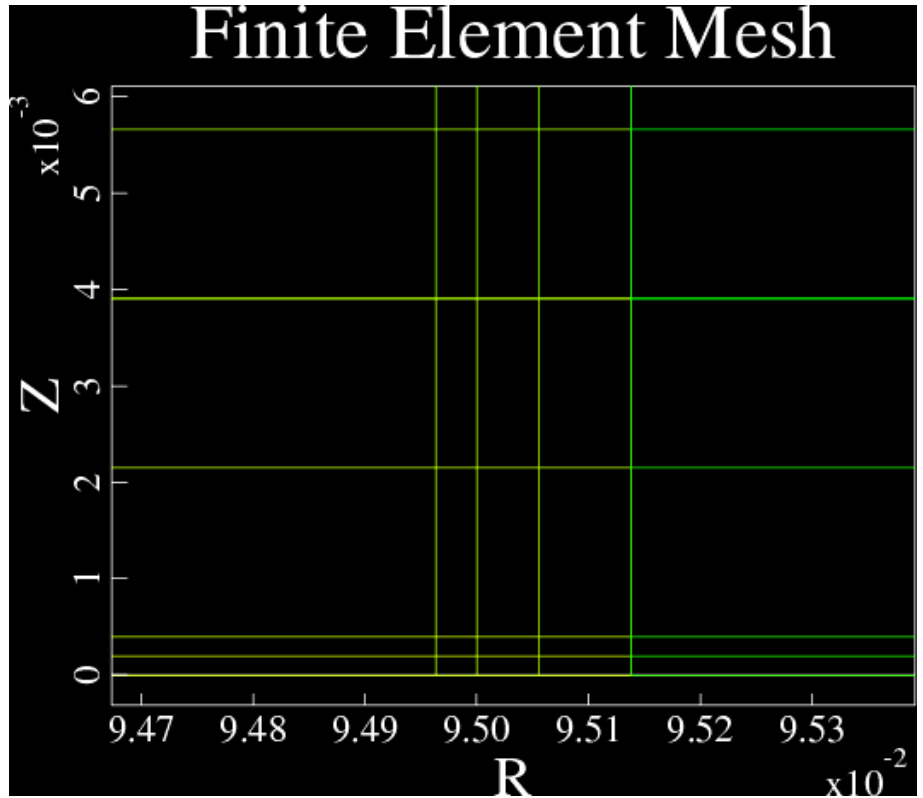
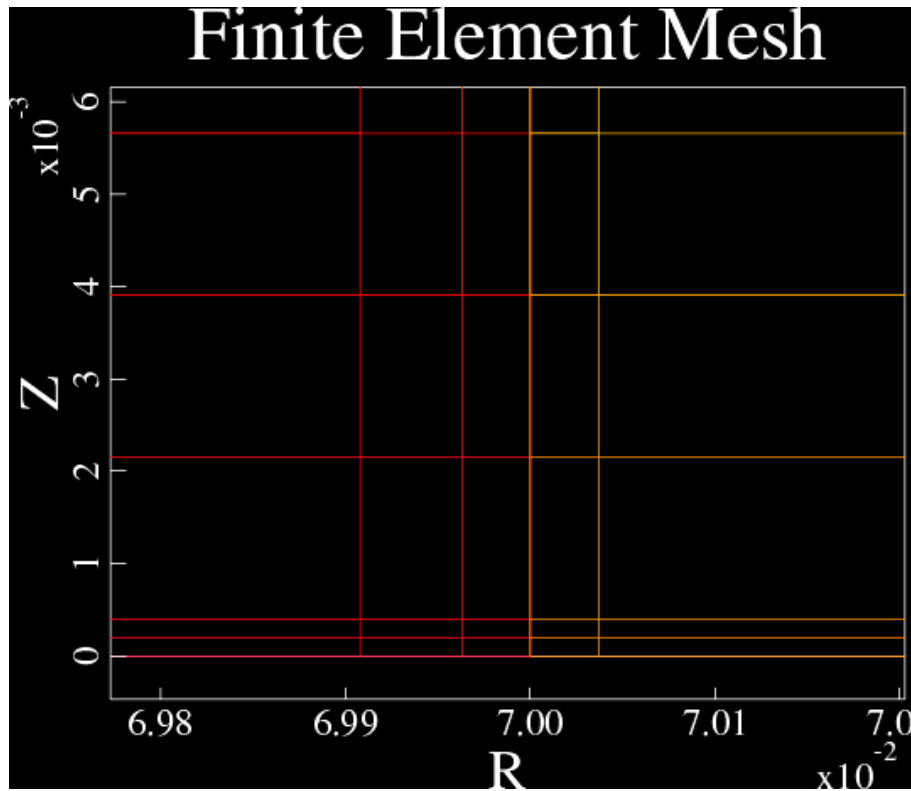
B_ϕ along $Z=0$ shows discontinuity



- 4 curves are at $t = [.53, .69, .82, .93] \mu s$
- note discontinuity at $R = .07$ and $.095$ where boundary condition is specified
- note strong effect of diffusive layer

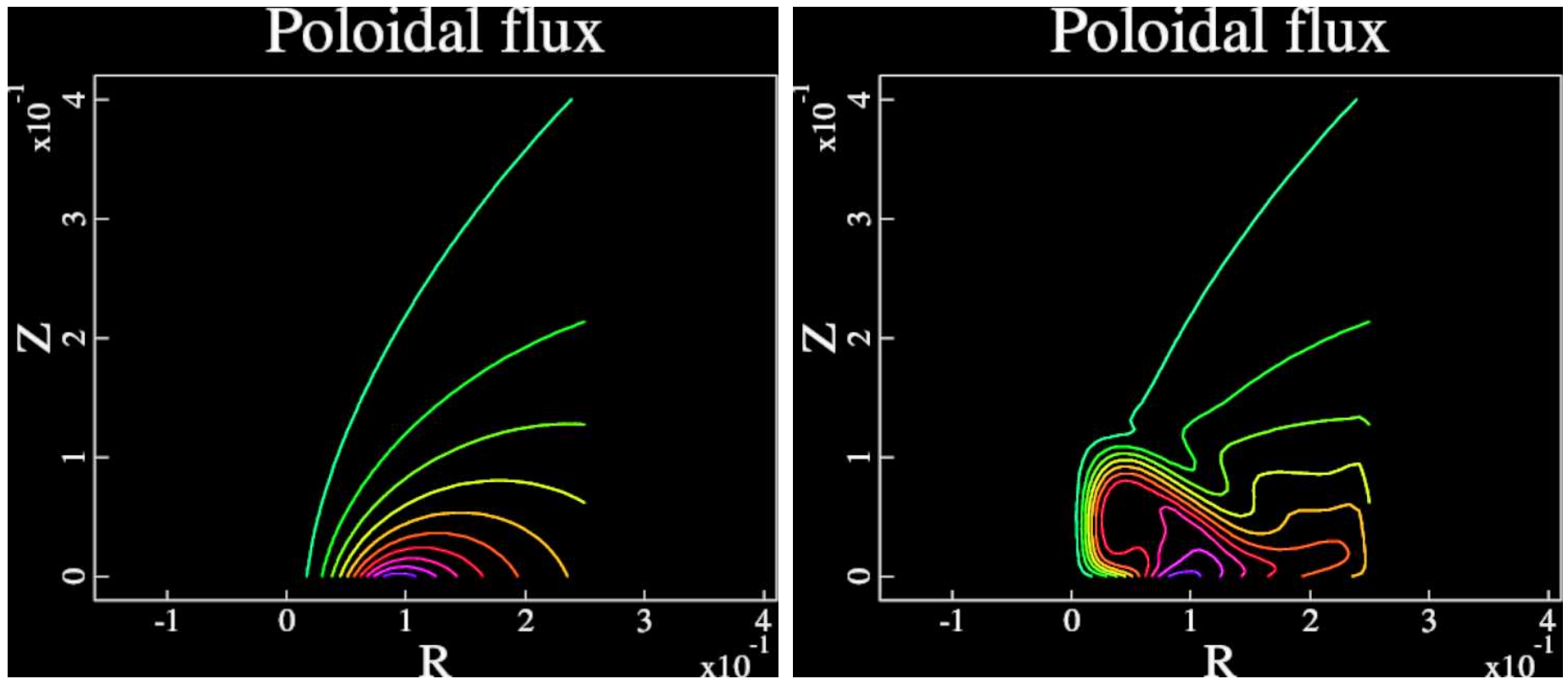
Zoomed in view of Coplanar Injection grid

discontinuous jump in $\frac{\partial B_\phi}{\partial R}$ requires high resolution



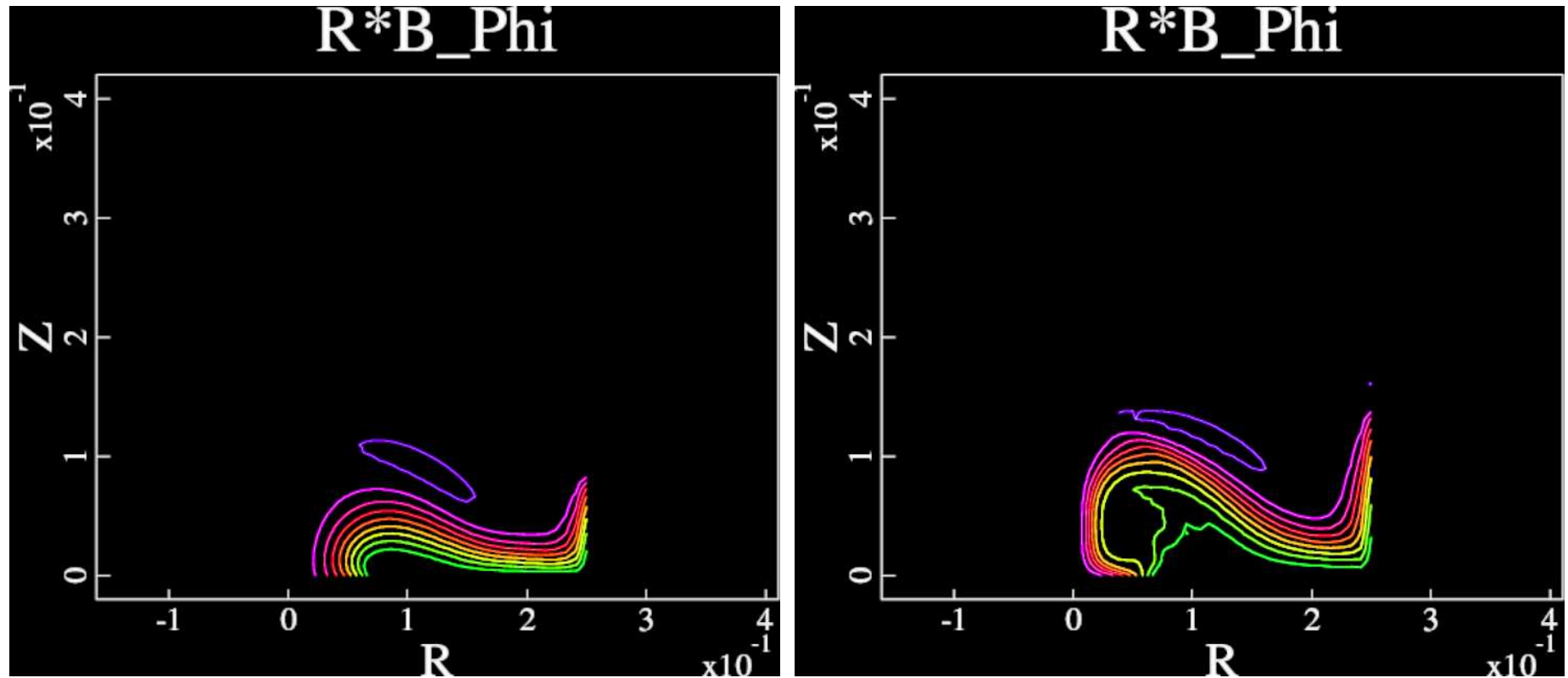
note the scale

Poloidal Flux Compression after $\sim 4.5\mu s$

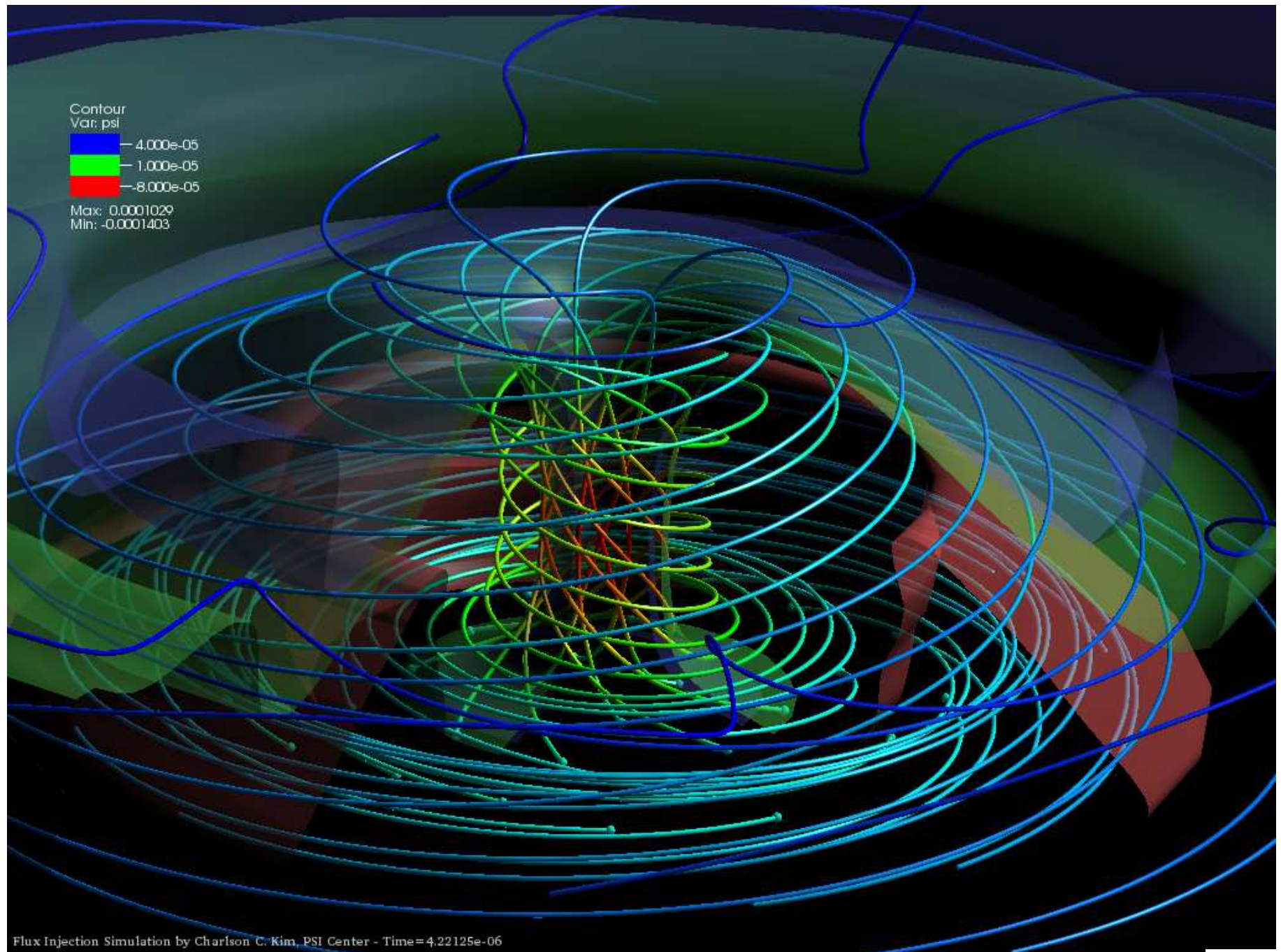


- 10 contour lines $[-.13, .9]mWb$
- contour lines same value in both plots

Evolution of Poloidal Current



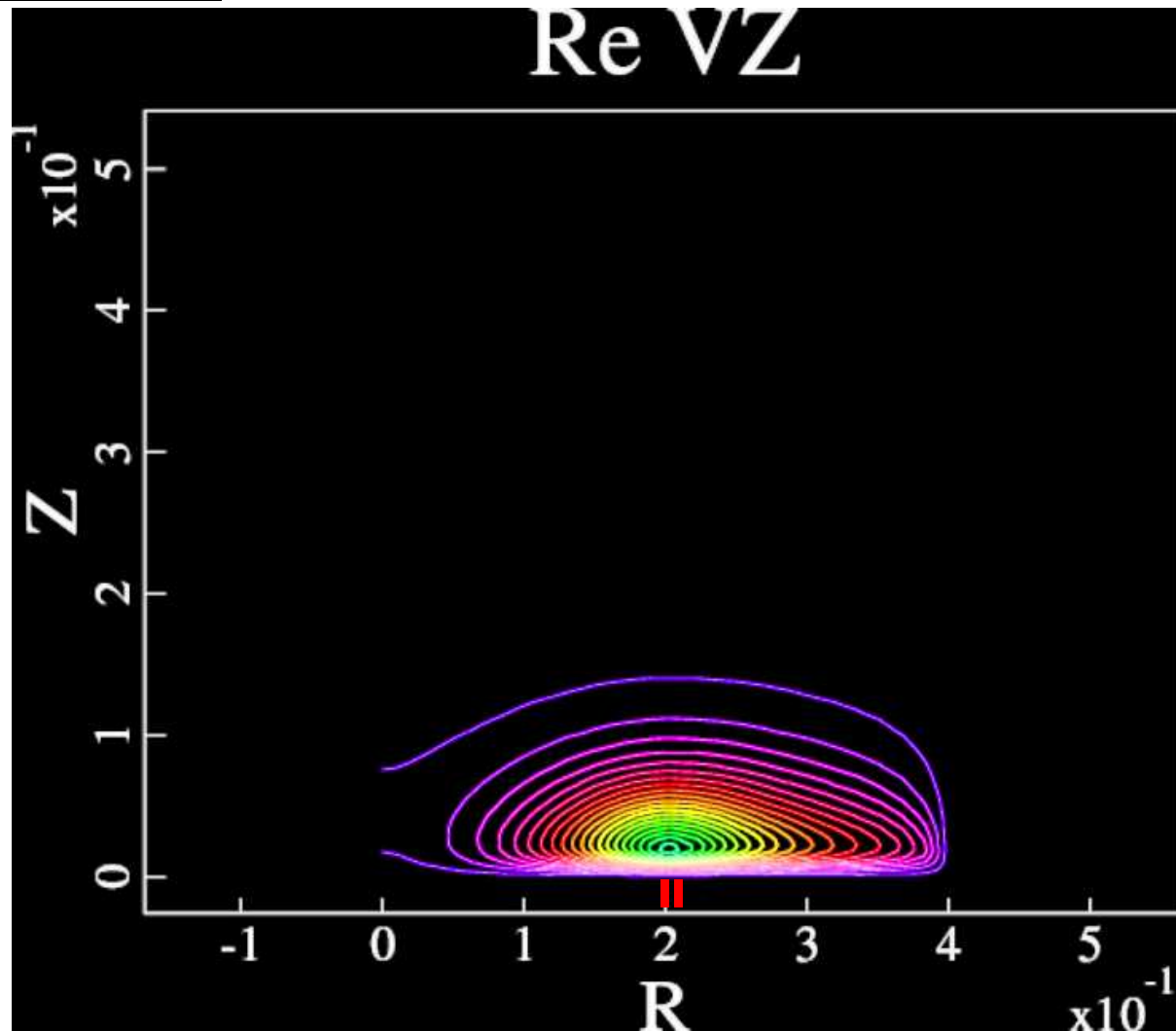
- left plot at $2.9\mu s$, contour range $[1.2, 11.7]kA$
- right plot at $4.2\mu s$, contour range $[1.8, 16.9]kA$
- current ramp time $10\mu s$, peak current $40kA$



Some Heuristic Considerations

- specifying B_ϕ meant to induce I_{pol}
- but fast injection also results in $\frac{\partial B_\phi}{\partial t} = -(\nabla \times \mathbf{E})_\phi$
- i.e. $\frac{\partial E_R}{\partial Z}$ and $\frac{\partial E_Z}{\partial R}$
- as B_ϕ increases \mathbf{E} becomes more E_\perp
- requires strong perpendicular flow in resistive model
- Hall physics removes inertial from field line evolution
- ? what about E_\parallel
 - in resistive model requires J_\parallel
 - Hall provides additional $\nabla_\parallel p_e$

Strong Alfvénic flows



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- as B_ϕ increases \mathbf{E} becomes more E_\perp
- requires strong perpendicular flow in resistive model
- Hall physics decouples inertia from field line motion
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 - in resistive model requires J_\parallel
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Summary, Ongoing Work, and Issues

- for resistive MHD fast injection results in strong Alfvénic flows
- Hall physics helps relieve strong flow issues
 - but could drive smaller scale fluctuations
 - still constrained by strong flows
 - are there parameters to fiddle
- flux gap model is incomplete
 - no-slip perfect conductor b.c. no longer applies
 - consistent b.c. remains to be resolved

Slides among friends



APS-DPP08 Dallas, TX



Initialize logical for setting boundary condition

```
IF (gridshape=='bell') THEN
```

```
  xov0=xo
```

```
  IF (wpack(1).gt.0._r8)
```

```
  $      xov0=xo-wpack(1)*ATANH(1._r8-2._r8*amp(1))*wpack(3)
```

```
  yov0=yo
```

```
  IF (wpack(2).gt.0._r8)
```

```
  $      yov0=yo+wpack(2)*ATANH(1._r8-2._r8*amp(2))*wpack(3)
```

C-----

C use xo and yo to set the current injection region

C-----

```
DO ibl=1,nrbl
```

```
  nv=seam(ibl)%nvert
```

```
  DO iv=1,nv
```

```
    ivn=iv-1
```

```
    IF (ivn==0) ivn=nv
```

```
    IF ((seam(ibl)%vertex(iv)%zgeom.EQ.ymin)
```

```
  $      .AND.(seam(ibl)%vertex(iv)%rgeom.GE.xov0)
```




```

$          .AND.(seam(ib1)%vertex(iv)%rgeom.LE.yov0)) THEN
          seam(ib1)%aplv0(1,iv)=.TRUE.
IF (seam(ib1)%excorner(iv))
$          seam(ib1)%aplv0(1,iv)=.FALSE.
IF ((seam(ib1)%vertex(ivn)%zgeom.EQ.ymin)
$          .AND.(seam(ib1)%vertex(ivn)%rgeom.GE.xov0)
$          .AND.(seam(ib1)%vertex(ivn)%rgeom.LE.yov0))
$          seam(ib1)%aplv0(2,iv)=.TRUE.
          write(*,*)'in bell',ib1,iv,'is true aplv0(2) is',
$                                     seam(ib1)%aplv0(2,iv)
          ENDIF
        ENDDO
      ENDDO
    ENDIF

```

Use the logical for applying boundary condition

```
DO ibl=1,SIZE(exblock_list)
  ibe=exblock_list(ibl)
  DO iv=1,seam(ibe)%nvert
    IF (seam(ibe)%applV0(1,iv)) THEN
      ix=seam(ibe)%vertex(iv)%intxy(1)
      iy=seam(ibe)%vertex(iv)%intxy(2)
      DO imode=1,nmodes
        IF (keff(imode)==0) be(ibe)%arr(3,ix,iy,imode)=
$          gun_cur/rb(ibe)%rz%fs(1,ix,iy)
$          *MAX(oneflg1,
$          .5_r8*(TANH((rb(ibe)%rz%fs(1,ix,iy)-xov0)/wpack(1))+1))
$          *MAX(oneflg2,
$          .5_r8*(TANH((yov0-rb(ibe)%rz%fs(1,ix,iy))/wpack(2))+1))
        ENDDO
      ENDIF
    ENDDO
  ENDDO
```