

Late Nonlinear Stage of Tokamak Ballooning Instability

The Effects of Hyperdiffusivity

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Can/how does ballooning filament break off from core plasma?

- ▶ Question:
 - ▶ What physics governs the late/terminal state of a ballooning filament?
- ▶ Previous work I: ELM-milestones (2005,2006)
 - ▶ DIII-D like equilibrium
 - ▶ Large diffusivity for numerical stability
- ▶ Previous work II: g-mode (2007)
 - ▶ 2D Shearless slab (purely interchange)
 - ▶ Low dissipation regime difficult to simulate
- ▶ This work: ballooning of ESC equilibrium
 - ▶ Toroidal coupling: energy cascades even linearly
 - ▶ 3D global structure: has to resolve parallel direction
 - ▶ Low dissipation regime: filaments carry high temperature
 - ▶ Physically intriguing and computationally challenging
 - ▶ What are the optimal physical and numerical dissipations?

A Solution: Hyperdiffusivity [Sovinec, 2008]

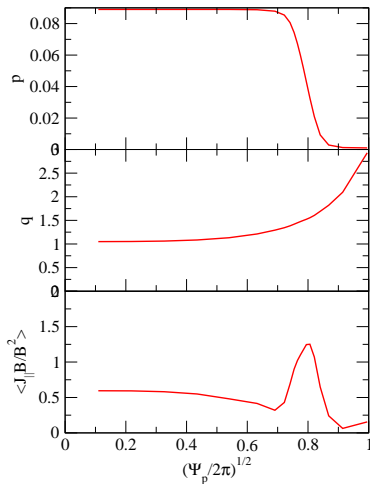
- ▶ Numerical difficulties mostly arise from cascading to small structures in density evolution
- ▶ Hyperdiffusivity improves numerical stability, while leaving macroscopic structures intact.

$$\partial_t n + \mathbf{u} \cdot \nabla n = -n \nabla \cdot \mathbf{u} - \nabla \cdot (D_h \nabla \nabla^2 n) \quad (1)$$

- ▶ Recently implemented hyperdiffusivity in `nimuw`
 - ▶ improves numerical stability of both linear and nonlinear runs;
 - ▶ allows physically ideal, dissipationless runs
 - ▶ simulations are able to reach later nonlinear stage of tokamak ballooning.

Tokamak equilibrium: Circular boundary and pedestal-like pressure

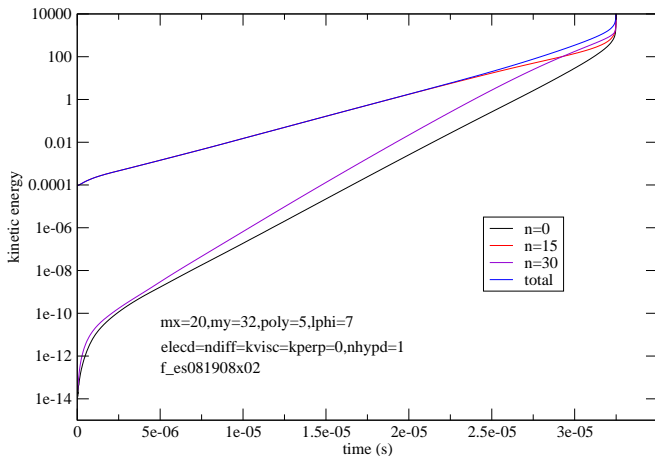
- ▶ $R_b = R_0 + a \cos(\theta + \delta \sin \theta)$,
 $Z_b = Z_0 + \kappa a \sin \theta$
- ▶ $Z_0 = 0$, $R_0 = 3$, $a = 1$,
 $\kappa = 1$, $\delta = 0$
- ▶ $q = q_0 [1 + (q_a/q_0 - 1)r^4]$
- ▶ $r = \sqrt{\Psi/\Psi_a}$, Ψ – toroidal flux
- ▶ $q_0 = 1.05$, $q_a = 3.0$, $B_0 = 1.0$,
 $\tau_A = R_0 q_0 / u_{A0} \sim 10^{-6} s$
- ▶ $p = p_p + h_p \tanh [(r_p - r)/L_p]$
- ▶ $p_p = 0.045$, $h_p = 0.044$,
 $r_p = 0.7$, $L_p = 0.05$



Hyperdiffusivity ensures a linear dissipationless ideal ballooning mode to run

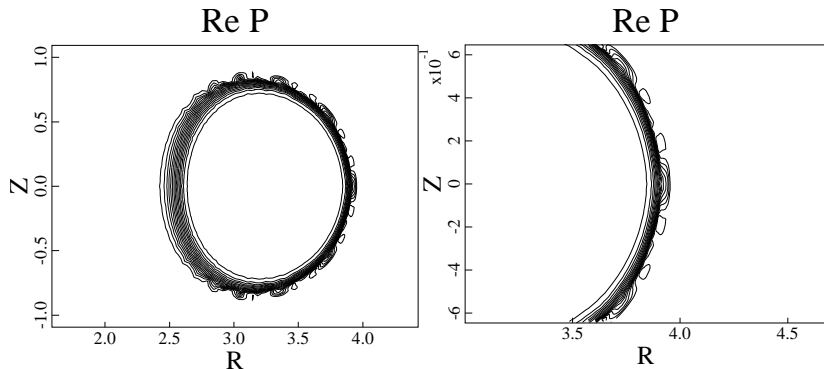
- ▶ $mx=20$, $my=32$, $poly=5$, $elec d=ndiff=kvisc=kperp=0$,
 $nd_hypd=1$
- ▶ The run was numerically unstable for $nd_hypd=0$.
- ▶ Growth rate insensitive to nd_hypd .

How far can a nonlinear ballooning simulation last (with hyperdiffusivity only)?



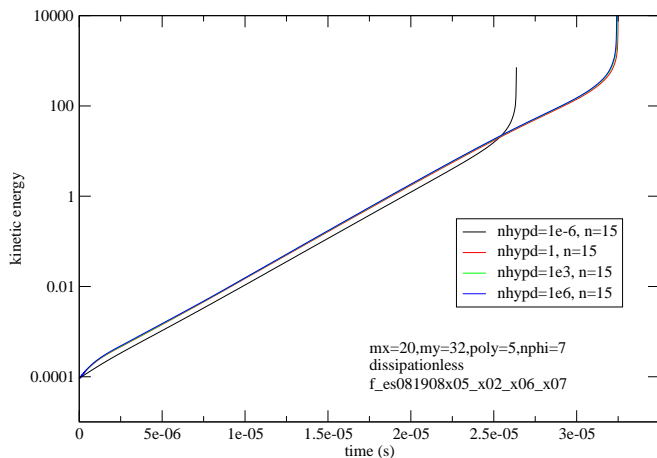
- ▶ I.C. n=15 nearly eigen-ballooning-mode
- ▶ Clearly nonlinear after $25\mu s$

Distortion of flux surface becomes visible near end of run

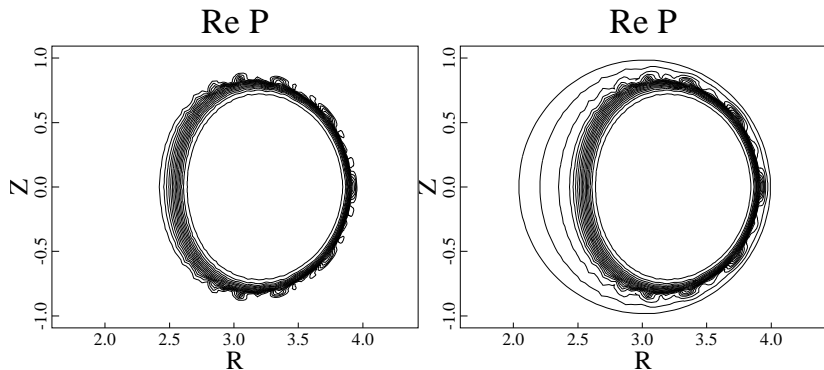


- ▶ Left: pressure contour at $30\mu s$; Right: zoom in of Left.
- ▶ $elec d=ndiff=kvisc=kperp=0$, $nhypd=1$.
- ▶ Displacement becomes comparable of pedestal width.
- ▶ $divB$ error large before numerical instability; code stops due to small dt .

Effects of hyperdiffusivity are mostly numerical



Distortion of flux surface appears better resolved with larger hyperdiffusivity



- ▶ Above: Pressure contours at $30\mu s$.
- ▶ Left: $n_{hypo} = 1e3$; Right: $n_{hypo} = 1e6$.
- ▶ $e_{elec} = n_{diff} = k_{visc} = k_{perp} = 0$.

Summary

- ▶ Continued study in late nonlinear stage of ballooning
- ▶ Recently developed hyperdiffusivity have been applied to the simulation.
- ▶ Initial results indicate improvement in numerical stability, extending the nonlinear stage of ballooning simulation in the low dissipation regime.
- ▶ Relevant physical dissipation regime to be identified.