

Implementations and applications of continuum closures

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Outline

- 1 Solve full plasma kinetic equation using continuum method.
 - Apply 2-D finite-elements (FE) in Chapman-Enskog-like kinetic equation.
 - Slowing down of diffuse ion beam.
- 2 Solve drift kinetic equation using continuum method.
 - Recap of recursive drift kinetic theory.
 - Application to sound wave damping problem.
 - One recursion introduces drift effects.

Solve Chapman-Enskog-like (CEL) kinetic equation.

- Couple solutions of 6-D, time-dependent Chapman-Enskog-like kinetic equation to evolving fluid equations through closures for \mathbf{q} and $\mathbf{\Pi}$:

$$\frac{\partial F}{\partial t} + \mathbf{v} \cdot \nabla F + \mathbf{a} \cdot \nabla F = C(F + f_M) - \frac{m}{T} (\mathbf{v}' \mathbf{v}' - v'^2 \frac{I}{3}) : \nabla \mathbf{u} - \frac{2f_m}{3\rho} L_1^{(3/2)} [\nabla \cdot \mathbf{q} + \mathbf{\Pi} : \nabla \mathbf{V} - Q] + \mathbf{v}' \cdot \left[\frac{f_m}{\rho} (\nabla \cdot \mathbf{\Pi} - \mathbf{R}) + \frac{f_m}{T} L_1^{(5/2)} \nabla T \right]$$

$$\text{where } f_M = n(\mathbf{x}, t) \left(\frac{m}{2\pi T(\mathbf{x}, T)} \right)^{3/2} e^{-\frac{m(\mathbf{v}-\mathbf{u})^2}{2T}}$$

$$\text{and } \mathbf{v} = v_{\parallel} \mathbf{b} + v_{\perp} (\mathbf{e}_2 \sin \gamma - \mathbf{e}_3 \cos \gamma).$$

Use 2-D FE, 1-D Fourier basis for F.

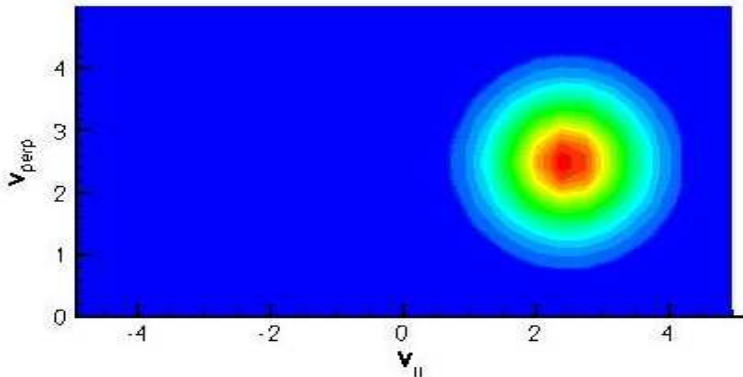
- Expand distribution function as
$$F = F_0 + \sum_{n>0} F_n e^{in\gamma} + F_n^\dagger e^{-in\gamma},$$
with Fourier coefficients expanded as
$$F_n = \sum_j F_{nj}(\mathbf{x}, t) \alpha_j(v_\perp, v_\parallel).$$
- At present, test particle portion of linearized Coulomb collision operator is implemented.

$$\int d\mathbf{v} \alpha C(f_\alpha, f_{M\beta}) =$$
$$-\frac{2\pi q_\alpha^2 q_\beta^2 \ln \Lambda_{ab}}{m_\alpha^2} \frac{n_\beta}{n_\alpha} \frac{1}{z_\beta} \int d\mathbf{v} \left(4 \frac{m_\alpha}{m_\beta} \frac{1}{v_{T\beta}} G(\mathbf{z}_\beta \cdot \frac{\partial \alpha}{\partial \mathbf{v}}) f_\alpha + \right.$$
$$\left. (E - G) \left(\frac{\partial \alpha}{\partial \mathbf{v}} \cdot \frac{\partial f_\alpha}{\partial \mathbf{v}} \right) + \left(\frac{3G - E}{z_\beta^2} \right) (\mathbf{z}_\beta \cdot \frac{\partial \alpha}{\partial \mathbf{v}}) (\mathbf{z}_\beta \cdot \frac{\partial f_\alpha}{\partial \mathbf{v}}) \right).$$

Solve full plasma kinetic equation using continuum method.
Solve drift kinetic equation using continuum method.

Apply 2-D finite-elements (FE) in Chapman-Enskog-like kinetic equation.
Slowing down of diffuse ion beam.

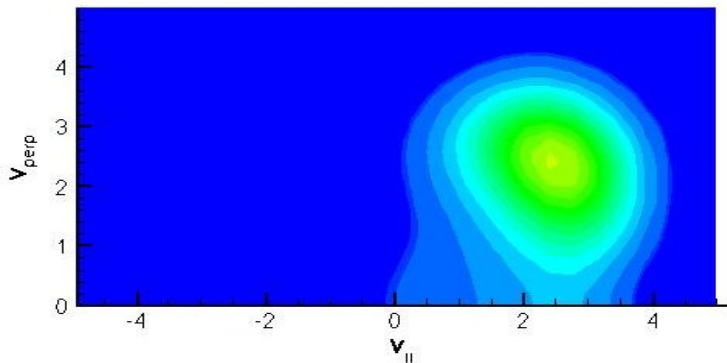
Apply test-particle operator to beam ions.



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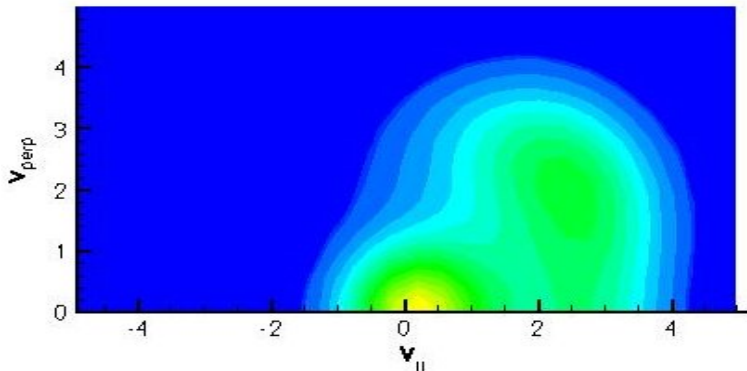
After 1 collision time.



Solve full plasma kinetic equation using continuum method.
Solve drift kinetic equation using continuum method.

Apply 2-D finite-elements (FE) in Chapman-Enskog-like kinetic equation.
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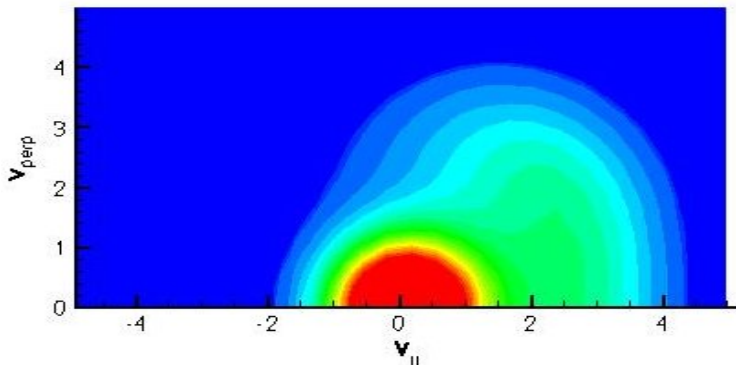
After 3 collision times.



Solve full plasma kinetic equation using continuum method.
Solve drift kinetic equation using continuum method.

Apply 2-D finite-elements (FE) in Chapman-Enskog-like kinetic equation.
Slowing down of diffuse ion beam.

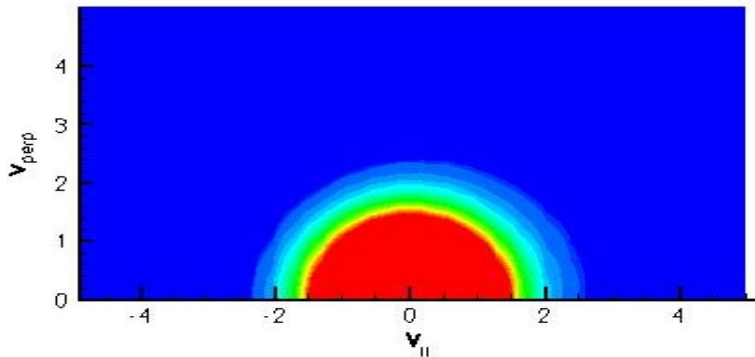
After 5 collision times.



Solve full plasma kinetic equation using continuum method.
Solve drift kinetic equation using continuum method.

Apply 2-D finite-elements (FE) in Chapman-Enskog-like kinetic equation.
Slowing down of diffuse ion beam.

After 20 collision times distribution is Maxwellian.



Future work on full continuum solution.

- Test-particle operator implemented.
- Implement field terms of linearized Coulomb operator.
- Implement nonlinear terms.
- Develop coding necessary to extend solution to spatial domain.

Drift kinetic theory a la Hazeltine '73.

- Define $L = d/dt - \Omega\partial/\partial\gamma$, total time derivative without the rapid gyro-motion. Here $\Omega = qB/m$.
- Hazeltine suggests solving
$$\Omega \frac{\partial \tilde{f}}{\partial \gamma} = -(L - C)f + \langle (L - C)f \rangle,$$
 by approximating f on right as \bar{f} (the gyro-angle independent part).
- Can repeat this by step by using $f = \bar{f} + \tilde{f}_1$ on right to get an improved f .
- Finally, insert $f = \bar{f} + \tilde{f}_j$ after j recursions into $\langle (L - C)f \rangle = 0$ and solve for \bar{f} .

Zero recursions ($\tilde{f} = 0$) leads to lowest-order CEL-drift kinetic equation.

- Solve lowest-order equation for electrons and ions to compute parallel closures, \mathbf{q}_{\parallel} and π_{\parallel} :

$$\frac{\partial F}{\partial t} + \mathbf{v}_{\parallel} \cdot \nabla F + \frac{q}{m} E_{\parallel} \frac{v_{\parallel}}{v} \frac{\partial F}{\partial v} =$$

$$\langle C(F + f_M) \rangle - \frac{mv^2}{T} f_M (\mathbf{b}\mathbf{b} - \frac{1}{3}) : \nabla \mathbf{u} -$$

$$+ \frac{2f_m}{3p} L_1^{(3/2)} \left[\nabla \cdot \mathbf{q} + \mathbf{\Pi} : \nabla \mathbf{V} - Q - S_0^{rf} \right] +$$

$$\mathbf{v}_{\parallel} \cdot \left[\frac{f_m}{p} (\nabla \cdot \mathbf{\Pi} - \mathbf{R} - F_0^{rf}) + \frac{f_m}{T} L_1^{(5/2)} \nabla T \right]$$

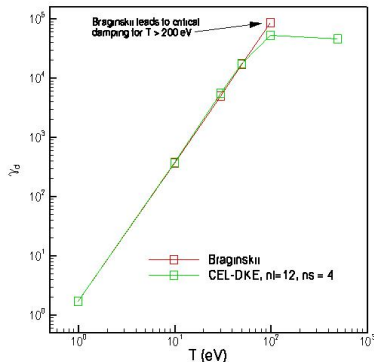
- Expand $F = \sum_{l=0}^{nl} F_l(\mathbf{x}, t, s) P_l(v_{\parallel}/v)$, where the coefficients $F_l(\mathbf{x}, t, s)$ are determined on a grid of ns grid points in the normalized speed, $s = v/v_T$.

Time-discretization and centering issues.

- Present implementation centers F_e and F_i with n , \mathbf{B} , T_e , T_i in time.
- n , \mathbf{B} , T_e and T_i dependence in Maxwellian drives approximately centered in implicit F advance.
- Advance \mathbf{V} , n , F_e, F_i , compute closures.
- Advance T_e, T_i and \mathbf{B} , solve again for F_e, F_i with approximately centered n , \mathbf{B} , T_e , T_i , and recompute closures.
- Advance T_e, T_i and \mathbf{B} with approximately centered closures.

Apply to sound wave damping problem.

- Test stress damping of sound waves with adiabatic equation of state.

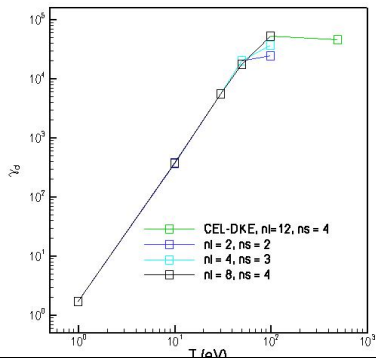


Solve full plasma kinetic equation using continuum method.
Solve drift kinetic equation using continuum method.

Recap of recursive drift kinetic theory.
Application to sound wave damping problem.
One recursion introduces drift effects.

Test convergence in velocity space.

- 1 Number of Legendre polynomials = n_l . Number of grid points in speed = n_s .
- 2 Limits on workstation with 4 GB of memory: $n_l=12$ and $n_s=4$.



Future work on lowest-order continuum solution to CEL-DKE.

- Implement RF quasilinear source terms imported from GENRAY.
- Implement finite flow corrections in Maxwellian drives.
- Implement calculation of self-consistent, equilibrium F_e and F_i (needed when there is equilibrium current and flow).
- Test continuum solution to CEL-DKE on various problems.

Compare with Chang/Callen '92.

- One-pole approximations for linear closures were given as:

$$\begin{aligned} \mathbf{b} \cdot \nabla \cdot \tilde{\mathbf{\Pi}} &\approx n_0 m_i \frac{3}{5} \left[\sqrt{\pi} k_{\parallel} v_{Ti} - \frac{6}{5} \left(\frac{\partial}{\partial t} + 0.32 \nu_{ii} \right) \right] \tilde{u}_{\parallel i} \\ &+ \frac{2}{5} \left[k_{\parallel} v_{Ti} - \sqrt{\frac{3\pi}{10}} \left(\frac{\partial}{\partial t} + 0.42 \nu_{ii} \right) \right] \frac{n_0}{k_{\parallel} v_{Ti}} (\nabla_{\parallel} \tilde{T}) \\ \tilde{q}_{\parallel} &\approx \frac{2}{5} \left[k_{\parallel} v_{Ti} - \sqrt{\frac{3\pi}{10}} \left(\frac{\partial}{\partial t} + 0.42 \nu_{ii} \right) \right] \frac{\rho_i}{k_{\parallel} v_{Ti}} \tilde{u}_{\parallel i} \\ &- \frac{9}{5\pi} \left[\sqrt{\pi} k_{\parallel} v_{Ti} - \frac{72}{25\pi} \left(\frac{\partial}{\partial t} + 0.42 \nu_{ii} \right) \right] \frac{n_0}{k_{\parallel}^2} (\nabla_{\parallel} \tilde{T}) \end{aligned}$$

- Compare linear damping results between theory, kinetic MHD and continuum closure calculations.

Drift corrections needed for many problems.

- Do one level of recursive theory to incorporate drift effects.
- Integrate, $\Omega \frac{\partial \tilde{f}_1}{\partial \gamma} = -(L - C)\bar{f} + \langle (L - C)\bar{f} \rangle$.
- Insert $f = \bar{f} + \tilde{f}_1$ into $\langle (L - C)f \rangle = 0$ and solve for \bar{f} .
- Solve: $\frac{\partial \bar{F}}{\partial t} + (\mathbf{v}_{\parallel} + \mathbf{v}_D) \cdot \nabla \bar{F} + a \frac{\partial \bar{F}}{\partial \epsilon} =$ Maxwellian drives including drift effects.