

# Hyper-resistivity and Other “Stabilization” Methods

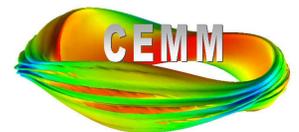
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# Thesis

Like other implicit differential operators of 4<sup>th</sup> order, hyper-dissipation can serve to “stabilize” the Hall B-advance.

## Outline

- Introduction
- Hyper-dissipation implementation
- Initial RFP-related results
- Discussion points

# Introduction

- The numerical stability analysis of the implicit leapfrog in (Sovinec and King, JCP **229**) focuses on the “semi-discrete” temporal advance, i.e. spatial approximation is not considered.
- Considering the function-space for the differential operators is also required with finite-elements.
- When the Hall term dominates the **B**-advance at all represented scales, that equation does not have the property of being mathematically coercive.

Wolfram MathWorld, “A bilinear functional  $\phi$  on a normed space  $E$  is called coercive (or elliptic) if there exists a positive constant such that

$$\phi(x, x) \geq K \|x\|^2$$

for all  $x \in E$  .

## The anisotropic, indefinite nature of the Hall term implies that it is not coercive.

- The cross product is anti-symmetric, and the dot product indicates a nontrivial null space (kernel).

$$\int \left[ \mathbf{A}^* \cdot \Delta \mathbf{B} - \frac{\Delta t}{2\mu_0 ne} \mathbf{B}^n \cdot (\nabla \times \mathbf{A}^* \times \nabla \times \Delta \mathbf{B}) \right] dVol = \text{stuff}$$

- A related point noted by Chacon, Lukin, and others is that resistive Hall-MHD without hyper-resistivity allows the 2fl reconnection scale to vanish.

$$E_{rec} \sim \frac{\eta}{\mu_0 \delta} B \sim \frac{k}{\mu_0 ne} B^2$$

There is no scale where dissipation dominates.

- Hyper-dissipation is one way to improve the properties of the Hall B-advance.

## Formulation: The NIMROD representation requires an auxiliary field for hyper-dissipation.

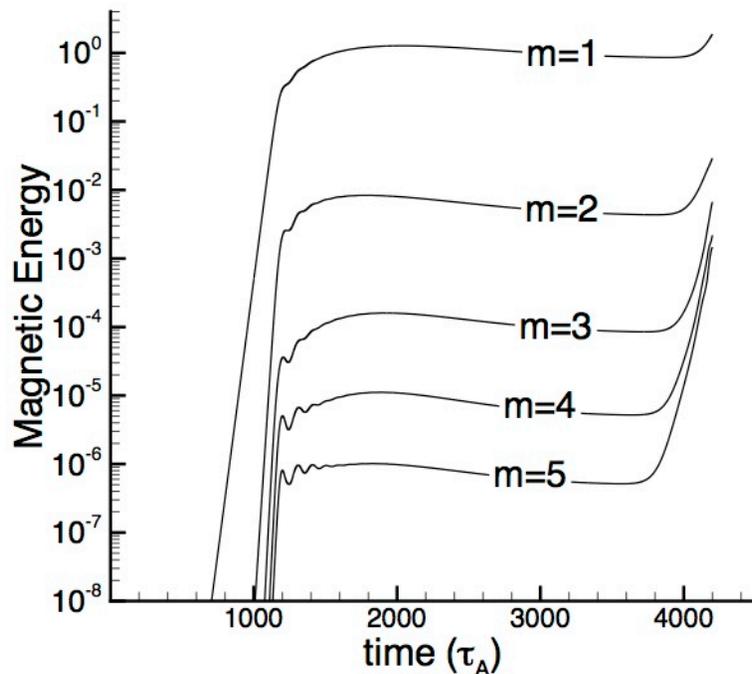
- For simplicity, consider just the Hall + hyper-dissipation terms, for all  $\mathbf{A}$  and  $\mathbf{F}$  in the vector-function space:

$$\int \left[ \mathbf{A}^* \cdot \Delta \mathbf{B} - \frac{\Delta t}{2\mu_0 ne} \mathbf{B}^n \cdot (\nabla \times \mathbf{A}^* \times \nabla \times \Delta \mathbf{B}) \right. \\ \left. - \sqrt{f\eta_h \Delta t} (\nabla \times \mathbf{A}^* \cdot \nabla \times \mathbf{G}) - \sqrt{f\chi_h \Delta t} (\nabla \cdot \mathbf{A}^* \nabla \cdot \mathbf{G}) \right. \\ \left. \mathbf{F}^* \cdot \mathbf{G} + \sqrt{f\eta_h \Delta t} (\nabla \times \mathbf{F}^* \cdot \nabla \times \Delta \mathbf{B}) + \sqrt{f\chi_h \Delta t} (\nabla \cdot \mathbf{F}^* \nabla \cdot \Delta \mathbf{B}) \right] dVol \\ = \int \left[ \frac{\Delta t}{\mu_0 ne} \mathbf{B}^n \cdot (\nabla \times \mathbf{A}^* \times \nabla \times \mathbf{B}^n) \right. \\ \left. - \sqrt{\frac{\eta_h \Delta t}{f}} (\nabla \times \mathbf{F}^* \cdot \nabla \times \mathbf{B}^n) - \sqrt{\frac{\chi_h \Delta t}{f}} (\nabla \cdot \mathbf{F}^* \nabla \cdot \mathbf{B}^n) \right] dVol$$

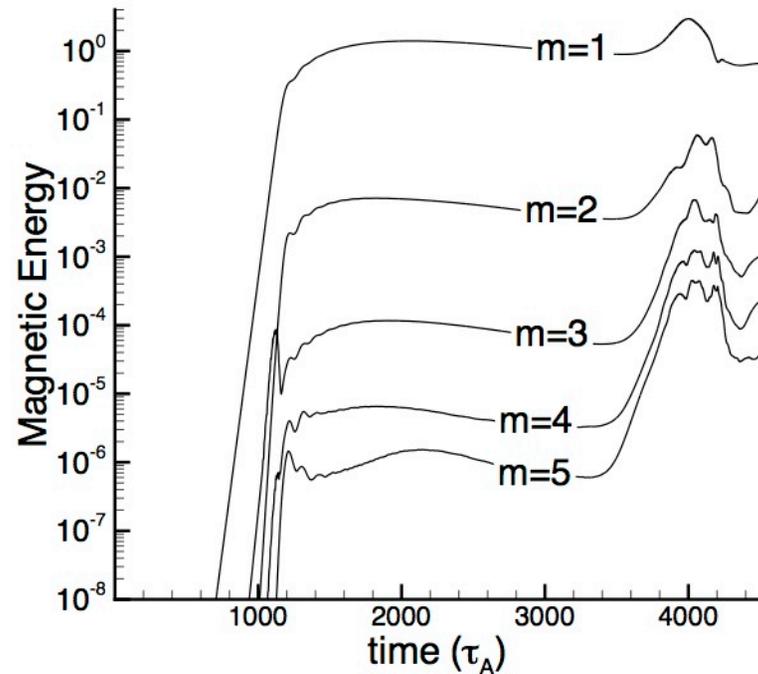
- Here,  $\eta_h$  is the hyper-magnetic diffusivity and  $\chi_h$  is a hyper-divergence-cleaning diffusivity.

# Results: The hyper-resistivity and hyper-cleaning operations have been applied in cylindrical two-fluid RFP computations.

- Josh has run S=20k computations with 120×64 meshes of degree 3.
- Results from a 60×32 mesh of degree 3 are compared below.



Fluctuation energy evolution w 120×64.



Fluctuation energy evolution w 60×32.

- Both computations have  $\eta / \mu_0 = 5 \times 10^{-5}$ ,  $\eta_h = 4 \times 10^{-8}$ , and Pm=1.
- The low-resolution case also has  $\chi_b = 1 \times 10^{-3}$ .

Computations with  $Pm=0.1$ , low resolution, and varied  $\eta_h$  indicate how the hyper-dissipation terms improve the numerical properties of the implicit operator.

- With  $\eta_h = 4 \times 10^{-9}$ , the 3D solver takes  $\sim 40$  iterations to converge the B-advance at small perturbation amplitude then fails after about 1000 steps.
- With  $\eta_h = 4 \times 10^{-8}$ , the 3D solver takes 7 iterations at small perturbation amplitude.
- With  $\eta_h = 4 \times 10^{-7}$ , the 3D solver takes 6 iterations at small perturbation amplitude.

## Discussion Points

- The hyper-dissipation term for the B-advance provide another tool to help 2-fluid computations.
- Their use will require care to check that these terms (like other numerical dissipation terms) do not affect converged results.
- The 6-vector solve is more costly than 3-vector solves, but the added robustness is expected to be worth the cost in some cases.
- Other 6-vector solves can be developed by modifying the new terms.