

Implementation of parallel moment equations in NIMROD

Hankyu Lee, Eric D. Held, Jeong-Young Ji

Utah State University

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Motivation

- Braginskii closures in Nimrod
 - In large Knudsen number, trying to capture some kinetic effects about parallel dynamics for magnetized plasmas
- Integral closures in Nimrod
 - Adiabatic and linear approximations to obtain the integral closures
 - Trying to incorporate non-adiabatic and nonlinear effects into parallel closures
 - Calculating exact parallel moment equations without integrals along the field lines

Total velocity moment equations

- Simpler than random velocity moments

- Landau (Fokker-Planck) kinetic equation

$$\frac{\partial f_a}{\partial t} + \mathbf{v} \cdot \nabla f_a + \frac{q_a}{m_a} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f_a}{\partial \mathbf{v}} = \sum_b C(f_a, f_b)$$

- Moment expansion

$$f_a(t, \mathbf{x}, \mathbf{v}) = f_a^M \sum_{lk} \mathbf{M}_a^{lk}(t, \mathbf{x}) \cdot \hat{\mathbf{P}}_a^{lk} \quad \text{with} \quad f_a^M = \frac{n_a}{\pi^{3/2} v_{Ta}^3} \exp(-s_a^2)$$

$$\mathbf{s}_a = \frac{\mathbf{v}_a}{v_{Ta}}, \mathbf{c}_a = \frac{\mathbf{v}_a - \mathbf{V}_a}{v_{Ta}} \quad (\text{in random velocity moments})$$

$$N_a^{lk} \equiv n_a \mathbf{M}_a^{lk}(t, \mathbf{x}) = \int d\mathbf{v} \hat{\mathbf{P}}_a^{lk} f_a(t, \mathbf{x}, \mathbf{v})$$

where $\hat{\mathbf{P}}_a^{lk}$'s are orthonormal polynomials of \mathbf{s}_a

Several low order moment equations

$P^{lk}(\mathbf{s}) = P^l(\mathbf{s})L_k^l(s^2)$	$L_k^l = L_k^{(l+\frac{1}{2})}$	M^{lk}	fluid moments	indep.
$P^0 = 1$	$L_0^0 = 1$	1	1	1
	$L_1^0 = \frac{3}{2} - s^2$	M^{01}	$-\sqrt{\frac{2}{3}} \frac{V^2}{v_T^2}$	1
$P^1 = \mathbf{s}$	$L_0^1 = 1$	M^{10}	$\sqrt{2} \frac{\mathbf{V}}{v_T}$	3
	$L_1^1 = \frac{5}{2} - s^2$	M^{11}	$-\sqrt{\frac{4}{5}} \frac{1}{nv_T T} (\mathbf{H} - \frac{5}{2} p \mathbf{V})$	3

When $\mathbf{V} \rightarrow 0$, $(\mathbf{H} - \frac{5}{2} p \mathbf{V}) \rightarrow \mathbf{h}$ in the random velocity moment expression

- Variables are $N_a^{lk} \equiv n_a M_a^{lk}(\mathbf{t}, \mathbf{x})$

Parallel moment equations in NIMROD

- Total velocity moment equations (no $\nabla \mathbf{V}$ coupled terms)[Ji, Jeong-Young, and Eric D. Held. "A framework for moment equations for magnetized plasmas." Physics of Plasmas 21.4 (2014): 042102.]

$$\begin{aligned} \partial_t \bar{N}_{\parallel}^{lp} + \sum_k \left[\bar{\Xi}_{pk}^l (\partial_t \ln T) \bar{N}_{\parallel}^{lk} + v_T (\bar{\Psi}_{pk}^{l+} \partial_{\parallel}^{l+} + \bar{\Phi}_{pk}^{l+} \partial_{\parallel} \ln T + \frac{q}{2T} \bar{\Theta}_{pk}^{l+} E_{\parallel}) \bar{N}_{\parallel}^{l+1,k} \right. \\ \left. + v_T (\bar{\Psi}_{pk}^{l-} \partial_{\parallel}^{l-} + \bar{\Phi}_{pk}^{l-} \partial_{\parallel} \ln T + \frac{q}{2T} \bar{\Theta}_{pk}^{l-} E_{\parallel}) \bar{N}_{\parallel}^{l-1,k} \right] = \sum C^{lpk} \bar{N}_{\parallel}^{lk} \\ + C^{(2)lp} (\bar{N}_{\parallel}, \bar{N}_{\parallel}) + G^{lp} \end{aligned}$$

$$\partial_{\parallel}^{l+} = \partial_{\parallel} - \frac{l+2}{2} (\partial_{\parallel} \ln B)$$

$$\partial_{\parallel}^{l-} = \partial_{\parallel} + \frac{l-1}{2} (\partial_{\parallel} \ln B)$$

- $(l, k) = (0, 0), (1, 0), (0, 1)$ driving terms are in the rhs.
- For example, $\frac{1}{n} \partial_{\parallel}^{2-} \bar{N}_{\parallel}^{1,0} = \partial_{\parallel}^{2-} \bar{M}_{\parallel}^{1,0} = (\partial_{\parallel} + \frac{1}{2} (\partial_{\parallel} \ln B)) (\sqrt{2} \frac{\mathbf{b} \cdot \mathbf{V}}{v_T}) = \mathbf{b} \cdot \nabla \sqrt{2} \frac{\mathbf{b} \cdot \mathbf{V}}{v_T} + \frac{1}{2} (\partial_{\parallel} \ln B) (\sqrt{2} \frac{\mathbf{b} \cdot \mathbf{V}}{v_T})$ (by using parallel strain tensor)

Time advance scheme

- In matrix form

$$\partial_t [n]^{\text{even}} = [c] [n]^{\text{even}} + [A] [n]^{\text{odd}} + \left[C^{(2)}(n, n) \right] + [g(n, T, \nabla \mathbf{V})]$$

$$\partial_t [n]^{\text{odd}} = [c] [n]^{\text{odd}} + [A] [n]^{\text{even}} + \left[C^{(2)}(n, n) \right] + [g(n, T, \mathbf{V}_{\text{ei}})]$$

$$[A] = -[\psi] \partial_{\parallel}^{\pm} - [\phi] \partial_{\parallel} \ln T - [\theta] E_{\parallel}$$

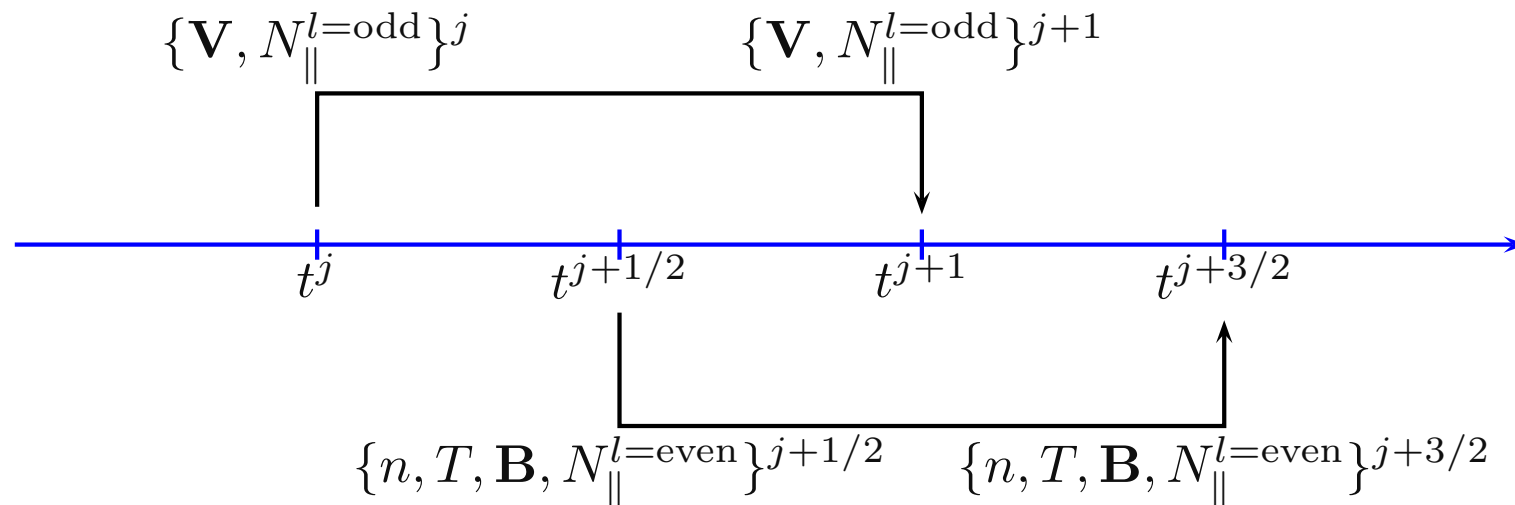
- Predictor-corrector time advance $t^k \rightarrow t^{k+1}$

$$(1 + \Delta t f_{[n]} [c]) [\Delta n]_{\text{pass}} = \Delta t ([c] [n]^k + [A] [n]^{k+1/2} + \left[C^{(2)}([n]^*, [n]^*) \right] + [g]^*)$$

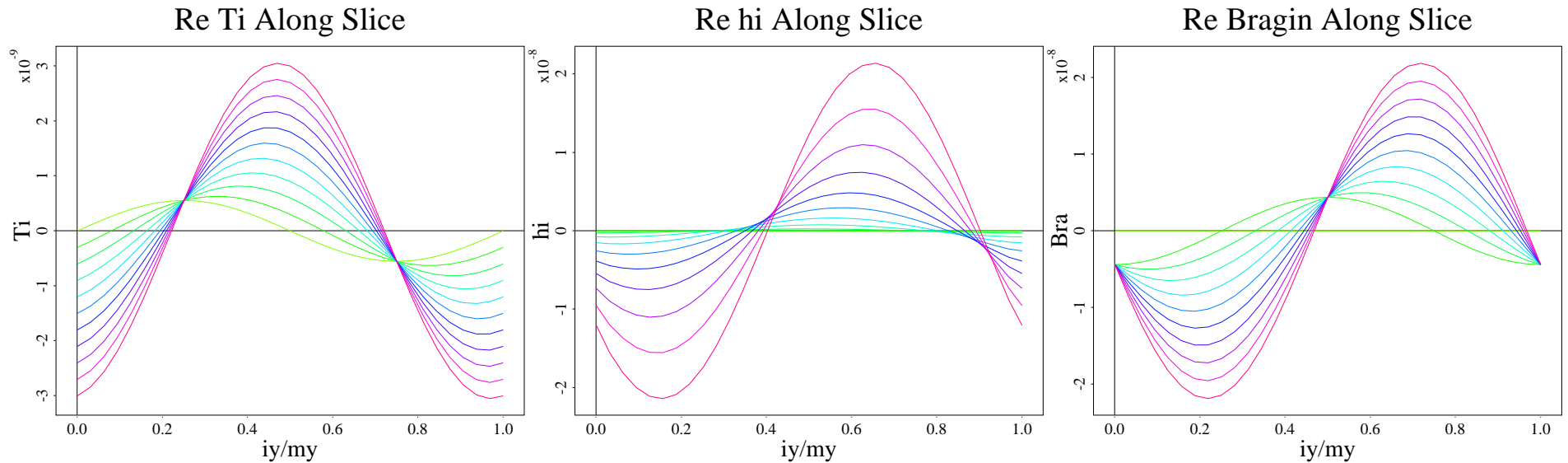
- For pass=predict, $[n]^* = [n]^k \rightarrow [\Delta n]_{\text{predict}}$
- For pass=correct, $[n]^{k+1} = [n]^k + f_{[n]} [\Delta n]_{\text{predict}}$
- Currently implemented for $f_{[n]} = 0$ and no $C^{(2)}$ term

Leap-frog scheme between even and odd moments

- NIMROD's flow velocity (\mathbf{V}) is staggered 1/2 step from number density (n), temperature (T), and magnetic field (\mathbf{B}) [Sovinec *et al*, International Sherwood Fusion Theory Conference 2008]

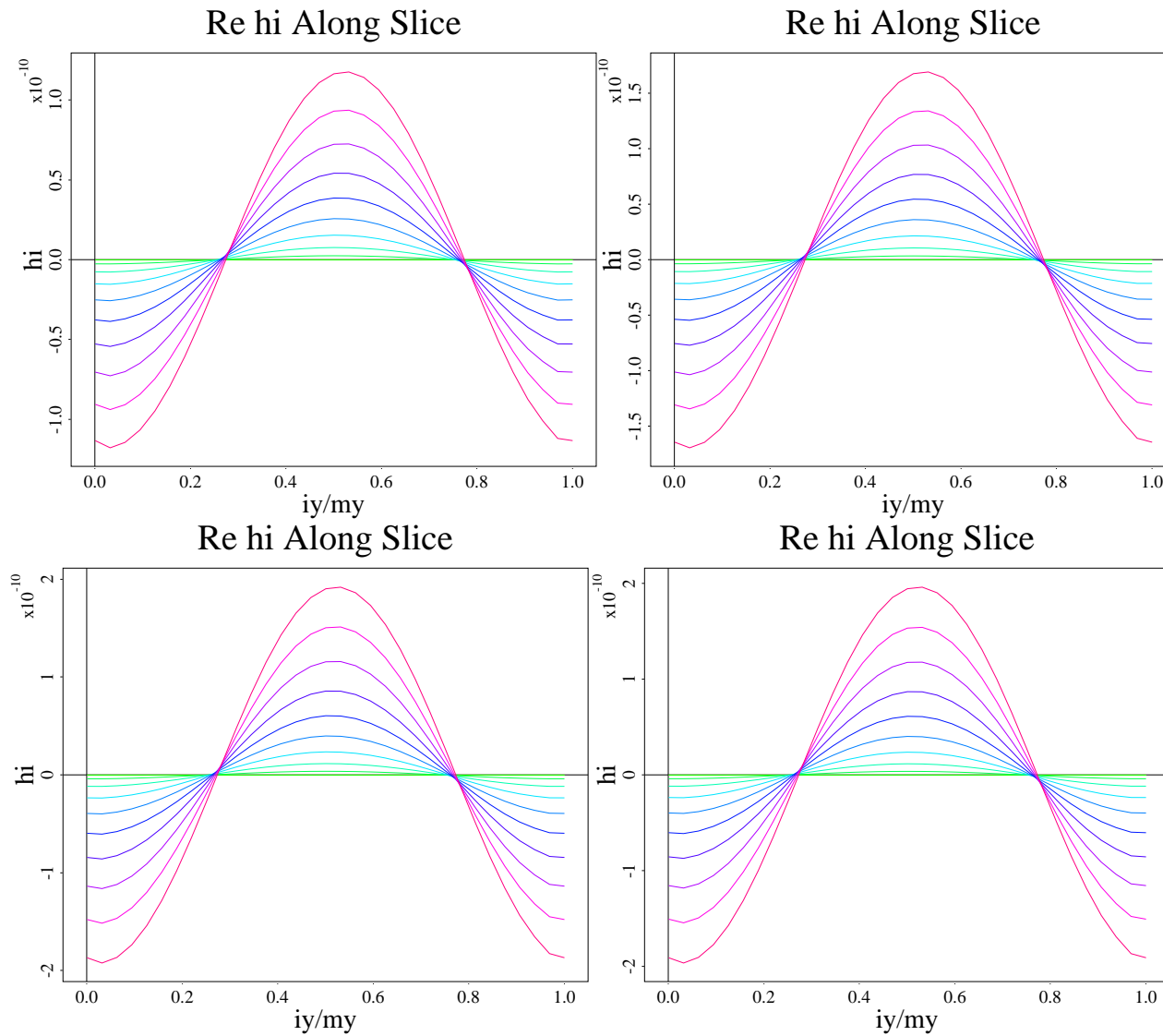


Verification by Ion acoustic wave simulation (ongoing)



- Ion temperature: 0.9 eV
- Using until $(l, k) = (2, 2)$ moments. Almost same results with Braginskii closure

Convergence along number of moments



$(l \text{ max}, k \text{ max}) = (2,2), (4,4), (8,8), (10,10)$

Future work

- Verification in equilibrium states [Held, E. D., et al. "Verification of continuum drift kinetic equation solvers in NIMROD." Physics of Plasmas 22.3 (2015)]
 - High aspect ratio equilibrium
 - High beta DIII-D like equilibrium
- Applying to nonlinear and time dependent phenomena
 - Magnetic reconnection
 - Plasma sheath
 - Waves