

The δf PIC-MHD model in NIMROD

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Outline

- hybrid kinetic-MHD model
- PIC and δf and the δf PIC-MHD model
- the details
 - drift kinetic equation and slowing down distribution
 - the implementation - **some code**
 - * ...
- example to illustrate diagnostics
- status and future directions



Hybrid Kinetic-MHD Model



The Hybrid Kinetic-MHD Equations^a

- in the **limit** $n_h \ll n_0$, $\beta_h \sim \beta_0$, **quasi neutrality**, only modification of MHD equations is addition of the **hot particle pressure tensor** in the momentum equation:

$$\rho \left(\frac{\partial \mathbf{U}}{\partial t} + \mathbf{U} \cdot \nabla \mathbf{U} \right) = \mathbf{J} \times \mathbf{B} - \nabla p_b - \nabla \cdot \underline{\mathbf{p}}_h$$

the subscripts b, h denote the bulk plasma and hot particles

- ρ, \mathbf{U} **for the entire plasma**, both bulk and hot particle
- the steady state equation

$$\mathbf{J}_0 \times \mathbf{B}_0 = \nabla p_0 = \nabla p_{b0} + \nabla p_{h0}$$

- p_{b0} **is scaled** to accomodate hot particles
- assumes equilibrium hot particle pressure is **isotropic**

^aC.Z.Cheng, 'A Kinetic MHD Model for Low Frequency Phenomena', *J. Geophys. Rev* **96**, 1991

Linearized Momentum equation and $\delta \underline{\mathbf{p}}_h$

- evolved momentum equation is ($\mathbf{U}_s = 0$)

$$\rho_s \frac{\partial \delta \mathbf{U}}{\partial t} = \mathbf{J}_s \times \delta \mathbf{B} + \delta \mathbf{J} \times \mathbf{B}_s - \nabla \cdot \delta \underline{\mathbf{p}}_b - \nabla \cdot \delta \underline{\mathbf{p}}_h$$

- assume CGL-like form $\delta \underline{\mathbf{p}}_h = \begin{pmatrix} \delta p_{\perp} & 0 & 0 \\ 0 & \delta p_{\perp} & 0 \\ 0 & 0 & \delta p_{\parallel} \end{pmatrix}$

- evaluate pressure moment at a position \mathbf{x} is

$$\begin{aligned} \delta p(\mathbf{x}) &= \int m(v - V_h)^2 \delta f(\mathbf{x}, \mathbf{v}) d^3 v \\ &= \sum_{i=1}^N \delta f_i m (v_i - V_h)^2 \delta^3(\mathbf{x} - \mathbf{x}_i) \end{aligned}$$

where sum is over the particles, δf_i is the perturbed phase space volume, m mass of the particle, and V_h is the velocity moment of the particles

PIC and δf and the δf PIC-MHD model



Overview of the PIC method

- in principle, $f(\mathbf{x}(t), \mathbf{v}(t))$ contains everything **but** intractable
- PIC is a Lagrangian simulation of phase space $f(\mathbf{x}, \mathbf{v}, t)$
- PIC approximation

$$f(\mathbf{x}, \mathbf{v}, t) \simeq \sum_{i=1}^N g_i(t) S(\mathbf{x} - \mathbf{x}_i(t)) \delta(\mathbf{v} - \mathbf{v}_i(t))$$

where the sum is over the **number of particles** N , i denotes the **particle index**, g_i is the **phase space volume** represented by the particle

- PIC evolves $f(\mathbf{x}_i(t), \mathbf{v}_i(t))$ by
 - **deposit** moment of $f(\mathbf{x}_i(t), \mathbf{v}_i(t))$ on grid using $S(\mathbf{x} - \mathbf{x}_i)$
 - **field solve** using moment
 - **advance** $[\mathbf{x}_i, \mathbf{v}_i]$ using solved fields
- all dynamics is in particle motion
- typically **PIC is noisy**, can't beat $1/\sqrt{N}$

The δf PIC method^a

- δf PIC **reduces the discrete particle noise** associated with conventional PIC
- Vlasov Equation

$$\frac{\partial f(\mathbf{z}, t)}{\partial t} + \dot{\mathbf{z}} \cdot \frac{\partial f(\mathbf{z}, t)}{\partial \mathbf{z}} = 0$$

- **split phase space** distribution into steady state and evolving perturbation

$$f = f_{eq}(\mathbf{z}) + \delta f(\mathbf{z}, t)$$

- apply PIC to $\delta f(\mathbf{z}, t)$
- particles advance along **characteristics** $\dot{\mathbf{z}}$
- substitute f in Vlasov Equation to get δf evolution equations

$$\delta \dot{f} = -\dot{\tilde{\mathbf{z}}} \cdot \frac{\partial f_{eq}}{\partial \mathbf{z}}$$

using $\dot{\mathbf{z}} = \dot{\mathbf{z}}_{eq} + \dot{\tilde{\mathbf{z}}}$ and $\dot{\mathbf{z}}_{eq} \cdot \frac{\partial f_{eq}}{\partial \mathbf{z}} = 0$

^aA. Y. Aydemir, "A unified MC interpretation of particle simulations...", *Physics of Plasmas*, **1**, 1994

The δf PIC-MHD model

- get fields from NIMROD
- advance particles and δf

$$\begin{aligned}\mathbf{z}_i^{n+1} &= \mathbf{z}_i^n + \dot{\mathbf{z}}(\mathbf{z}_i)\Delta t \\ \delta f_i^{n+1} &= \delta f_i^n + \dot{\delta f}(\mathbf{z}_i)\Delta t\end{aligned}$$

- deposit moment

$$\delta p(\mathbf{x}) = \sum_{i=1}^N \delta f_i m (v_i - V_h)^2 S(\mathbf{x} - \mathbf{x}_i)$$

- advance NIMROD equations with the hybrid Kinetic-MHD momentum equation

$$\rho_s \frac{\partial \delta \mathbf{U}}{\partial t} = \mathbf{J}_s \times \delta \mathbf{B} + \delta \mathbf{J} \times \mathbf{B}_s - \nabla \cdot \delta \underline{\mathbf{p}}_b - \nabla \cdot \delta \underline{\mathbf{p}}_h$$



The Details



Drift Kinetic Equation of Motion

- the **drift kinetic** equations of motion are used as the particle **characteristics**

$$\dot{\mathbf{x}} = v_{\parallel} \hat{\mathbf{b}} + \frac{m}{eB^4} \left(v_{\parallel}^2 + \frac{v_{\perp}^2}{2} \right) \left(\mathbf{B} \times \nabla \frac{B^2}{2} \right) + \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \frac{\mu_0 m v_{\parallel}^2}{eB^2} \mathbf{J}_{\perp}$$

$$m\dot{v}_{\parallel} = -\hat{\mathbf{b}} \cdot (\mu \nabla B - e\mathbf{E})$$

- reduces $6D$ phase space variables to $4D + 1$ $\left[\mathbf{x}(t), v_{\parallel}(t), \mu = \frac{\frac{1}{2} m v_{\perp}^2}{\|\mathbf{B}\|} \right]$
- follows gyrocenter in the limit of **zero Larmor radius**

Slowing Down Distribution for Hot Particles

- the slowing down distribution function

$$f_{eq} = \frac{P_0 \exp\left(\frac{P_\zeta}{\psi_0}\right)}{\varepsilon^{3/2} + \varepsilon_c^{3/2}}$$

$P_\zeta = g\rho_{\parallel} - \psi$ is the **canonical toroidal momentum** and ε is the **energy**, ψ_0 is the **gradient scale length**, and ε_c is the **critical slowing down energy**

- the evolution equation for δf

$$\delta \dot{f} = f_{eq} \left\{ \frac{mg}{e\psi_0 B^3} \left[\left(v_{\parallel}^2 + \frac{v_{\perp}^2}{2} \right) \delta \mathbf{B} \cdot \nabla B - \mu_0 v_{\parallel} \mathbf{J} \cdot \mathbf{E} \right] + \frac{\delta \mathbf{v} \cdot \nabla \psi_p}{\psi_0} + \frac{3}{2} \frac{e\varepsilon^{1/2}}{\varepsilon^{3/2} + \varepsilon_c^{3/2}} \mathbf{v}_D \cdot \mathbf{E} \right\}$$

$$\mathbf{v}_D = \frac{m}{eB^3} \left(v_{\parallel}^2 + \frac{v_{\perp}^2}{2} \right) (\mathbf{B} \times \nabla B) + \frac{\mu_0 m v_{\parallel}^2}{eB^2} \mathbf{J}_{\perp}$$

$$\delta \mathbf{v} = \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \mathbf{v}_{\parallel} \cdot \frac{\delta \mathbf{B}}{B}$$

Summary of Equations

- kinetic-MHD momentum equation

$$\rho_s \frac{\partial \delta \mathbf{U}}{\partial t} = \mathbf{J}_s \times \delta \mathbf{B} + \delta \mathbf{J} \times \mathbf{B}_s - \nabla \cdot \delta \underline{\mathbf{p}}_b - \nabla \cdot \delta \underline{\mathbf{p}}_h \quad (1)$$

- pressure moment

$$\begin{pmatrix} \delta p_{\perp} \\ \delta p_{\parallel} \end{pmatrix} = \sum_{i=1}^N \delta f_i \begin{pmatrix} \mu \|\mathbf{B}\| \\ \frac{1}{2} m v_{\parallel}^2 \end{pmatrix} S(\mathbf{x} - \mathbf{x}_i) \quad (2)$$

- drift kinetic characteristics

$$\dot{\mathbf{x}} = v_{\parallel} \hat{\mathbf{b}} + \frac{m}{eB^4} \left(v_{\parallel}^2 + \frac{v_{\perp}^2}{2} \right) \left(\mathbf{B} \times \nabla \frac{B^2}{2} \right) + \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \frac{\mu_0 m v_{\parallel}^2}{eB^2} \mathbf{J}_{\perp} \quad (3)$$

$$m \dot{v}_{\parallel} = -\hat{\mathbf{b}} \cdot (\mu \nabla B - e \mathbf{E}) \quad (4)$$

- δf evolution equation

$$\delta \dot{f} = f_{eq} \left\{ \frac{mg}{e\psi_0 B^3} \left[\left(v_{\parallel}^2 + \frac{v_{\perp}^2}{2} \right) \delta \mathbf{B} \cdot \nabla B - \mu_0 v_{\parallel} \mathbf{J} \cdot \mathbf{E} \right] + \frac{\delta \mathbf{v} \cdot \nabla \psi_p}{\psi_0} + \frac{3}{2} \frac{e\epsilon^{1/2}}{\epsilon^{3/2} + \epsilon_c^{3/2}} \mathbf{v}_D \cdot \mathbf{E} \right\} \quad (5)$$

$$\mathbf{v}_D = \frac{m}{eB^3} \left(v_{\parallel}^2 + \frac{v_{\perp}^2}{2} \right) (\mathbf{B} \times \nabla B) + \frac{\mu_0 m v_{\parallel}^2}{eB^2} \mathbf{J}_{\perp} \quad (6)$$

$$\delta \mathbf{v} = \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \mathbf{v}_{\parallel} \cdot \frac{\delta \mathbf{B}}{B} \quad (7)$$

Implementation Details

- particle data structure
- initialization - importance sampling
 - spatial load
 - velocity load
- PIC in FEM
 - inverse mapping and the search/sort and parallelization
 - particle field evaluation
 - particle deposition

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Particle Data Structure part_type_mod.f

- particle data type

```
TYPE :: particle
  REAL(r8), DIMENSION(0:2) :: R,Z,P,v1,v2,v3,w
  REAL(r8) :: g0(0:1),NS(0:3),GNS(1:2,0:3),eta,xi
  INTEGER(i4) :: ix,iy,id,pass
END TYPE particle
```

- species data type composed of array of types

```
TYPE :: species
  CHARACTER(10) :: part_type,hot_distro,pnorm_field
  LOGICAL :: add_phot,phdump,pdbg,trace
  INTEGER(i4) :: number,ntot,MPI_PART,packsz
  REAL(r8) :: mass,charge,vhmx,ecrit,avu0,avw0,orbits
  REAL(r8) :: P0,R0,nh0,rh0,hwdth,pres0,dmphi,anisop
  REAL(r8) :: eparallel,wwgt(0:2),passwgt(0:1)
  TYPE(particle), DIMENSION(:), POINTER :: part
END TYPE species
```

Particle Type

- $6D$ phase space variables and weight

REAL(r8), DIMENSION(0:2) :: R,Z,P,v1,v2,v3,w

– for part_type='drift'

– v1 is v_{\parallel}

– v2(1:2) is $(\mu, \|\mathbf{B}\|^{-1})$

- REAL(r8) :: g0(0:1),NS(0:3),GNS(1:2,0:3),eta,xi

– g0 phase space volume

– NS,GNS shape function value and derivatives (for linear basis)

– eta,xi logical coordinate

- INTEGER(i4) :: ix,iy,id,pass

– ix,iy logical index

– id unique particle I.D.

– pass passed or trapped status



Species Type

- species descriptors and array of particles

```
TYPE :: species
```

```
  CHARACTER(10) :: part_type,hot_distro,pnorm_field
```

```
  LOGICAL :: add_phot,phdump,pdbg,trace
```

```
  INTEGER(i4) :: number,ntot,MPI_PART,packsz
```

```
  REAL(r8) :: mass,charge,vhmx,ecrit,avu0,avw0,orbits
```

```
  REAL(r8) :: P0,R0,nh0,rh0,hwdth,pres0,dmphi,anisop
```

```
  REAL(r8) :: eparallel,wwgt(0:2),passwgt(0:1)
```

```
  TYPE(particle), DIMENSION(:), POINTER :: part
```

```
END TYPE species
```

- CHARACTER(10) :: part_type,hot_distro,pnorm_field
 - **part_type**=[drift, boris] equations of motion
 - **hot_distro**=[monoenergetic_perp, monoenergetic, maxwellian, slowing down] velocity distribution function
 - **pnorm_field**=[pres_eq, uniformp, uniformn, gauss] spatial distribution



- LOGICAL :: add_phot, phdump, pdbg, trace
- INTEGER(i4) :: number, ntot, MPI_PART, packsiz
 - number local particle count
 - ntot global particle count
 - MPI_PART, packsiz mpi parameters
- REAL(r8) :: mass, charge, vhm_x, ecrit, avu0, avw0, orbits
 - vhm_x parameter for velocity distribution
 - ecrit ε_c slowing down critical energy in KeV
 - avu0, avw0 equilibrium offset (*unused*)
 - orbits number of subcycled particle steps
- REAL(r8) :: P0, R0, nh0, rh0, hwdth, pres0, dmphi, anisop
 - P0 ψ_0 in slowing down distribution Ppsi0 in nimrod.in
 - R0 major radius (*unused*)
 - nh0, rh0, hwdth $=nh0 \exp\left(-\frac{(r - rh0)^2}{(L * hwdth)^2}\right)$ for pnorm_field=gauss
 - pres0(*unused*)

- `dmphi` length in periodic direction
- `anisop` multiplier for anisotropic pressure, sometimes set to zero
- `REAL(r8) :: eparallel, wwt(0:2), passwt(0:1)`
 - `eparallel` multiplier for eparallel nonlinearity
 - `wwt(0:2)` multiplier for different portions of weight equation controlled by `weight_use=[all, velocity, spatial, gradn]`
 - `passwt(0:1)` multiplier for moment calculation for passed and trapped particles controlled by `phase_use=[all, pass, trap]`

Other parameters in &particle_input

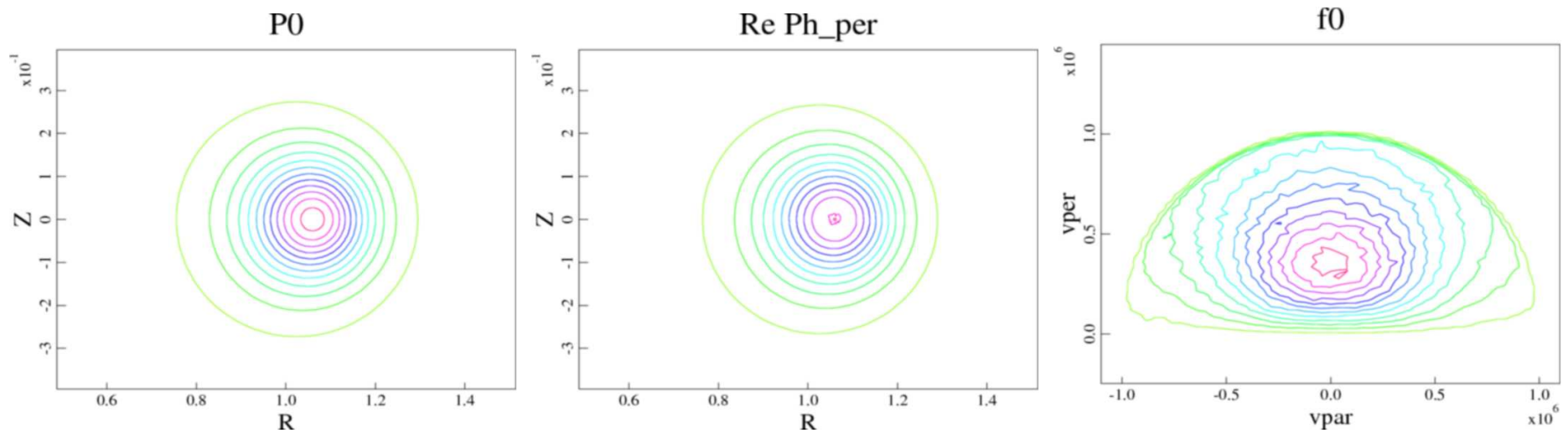
- **nm** total number of **requested** particles
- restart logical to read in `part_dump*.bin` file
- `phqty=2,6` number of components of pressure tensor
- **betafrac**=[0,1] fraction of pressure carried by hot particles
- `ph_pd` polynomial order for particles and phot
- `tpartdr` directory for `part_dump*.bin` file
- `bmxs,emxs` offset in `mx` to load annulus of particles
- `pdbg` turn on verbose output

Implementation Details

- particle data structure
- **initialization** - importance sampling
 - spatial load
 - velocity load
- PIC in FEM
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Initialization - mimic $f_{eq}(\mathbf{z})$

- particle **loading mimics $f_{eq}(\mathbf{z})$** in space and velocity - importance sampling



- example uses `pnorm_field=pres_eq`, `hot_distro=slowing down`
- `phot` in dumpfile if `phdump=.true.`
- diagnostic **equilibrium phot** is in `dump.99998,dump.99999`
- volume average $f(v_{\parallel}, v_{\perp}), \delta f(v_{\parallel}, v_{\perp})$ from `$>xd phase`

Algorithmic Details

- **nonuniform loading** in physical and velocity space
- match numeric load with analytic distribution $f_{eq}(\mathbf{x}, \mathbf{v})$
- physical space loaded proportional to **pnorm_field**
 - area weighted norm field is calculated on FE grid
 - particles are loaded cell by cell $\propto \frac{\Delta P_a}{P_a}$
- velocity space loaded to match **hot_distro**
 - ensure **isotropy** in \mathbf{v}
 - **first** initialize energy
 - compute uniform mapping function

$$D(\varepsilon) = \frac{\int_0^\varepsilon f(\varepsilon') d\varepsilon'}{\int_0^{\varepsilon^{max}} f(\varepsilon') d\varepsilon'}$$

- $\varepsilon = D^{-1}(\text{UniformRandomNumber})$
- **then** pick direction and project energy to velocity
- ** **$g\rho_{\parallel}$ not** accounted for in slowing down distribution



Implementation Details

- particle data structure
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PIC in FEM - nontrivial

- particles pushed in (R, Z) **but** field quantities (including (R, Z)) in (η, ξ)
- FE representation for (R, Z) needs to be **inverted**

$$R = \sum_j R_j N_j(\eta, \xi), \quad Z = \sum_j Z_j N_j(\eta, \xi),$$

- Use **Newton method** to solve for (η, ξ) given (R, Z)

$$\begin{Bmatrix} \eta^{k+1} \\ \xi^{k+1} \end{Bmatrix}_j = \begin{Bmatrix} \eta^k \\ \xi^k \end{Bmatrix}_j + \frac{1}{\Delta_j^k} \begin{pmatrix} \frac{\partial Z}{\partial \xi} & -\frac{\partial R}{\partial \xi} \\ -\frac{\partial Z}{\partial \eta} & \frac{\partial R}{\partial \eta} \end{pmatrix}_j^k \begin{Bmatrix} R - R^k \\ Z - Z^k \end{Bmatrix}_j$$

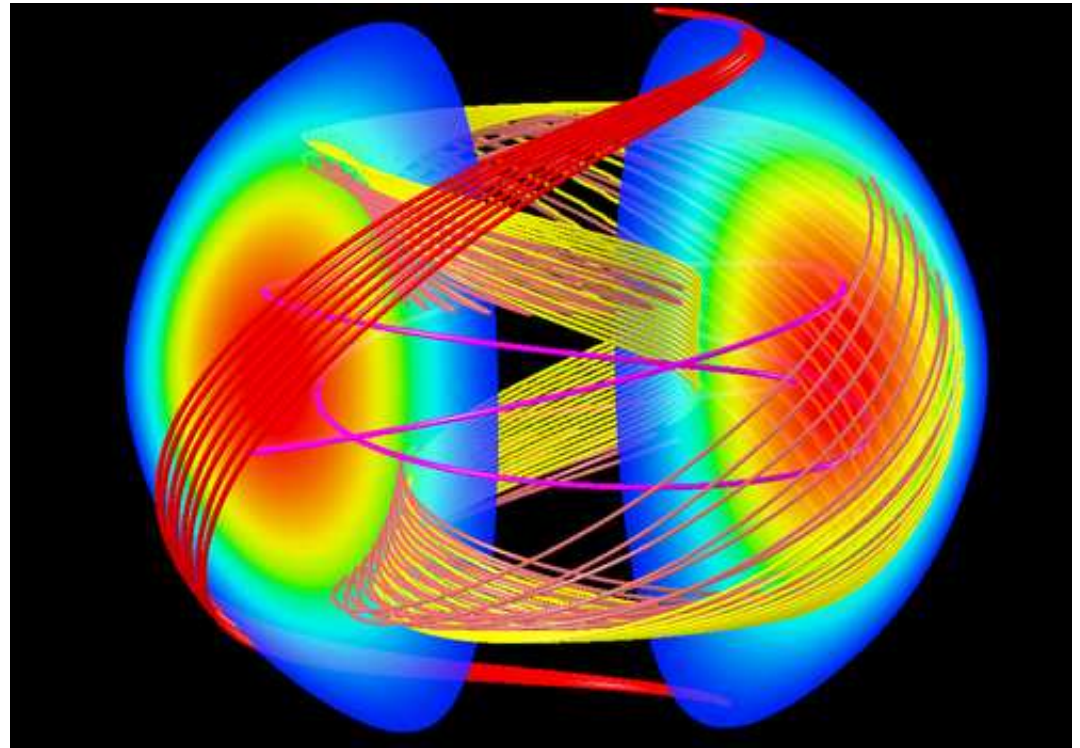
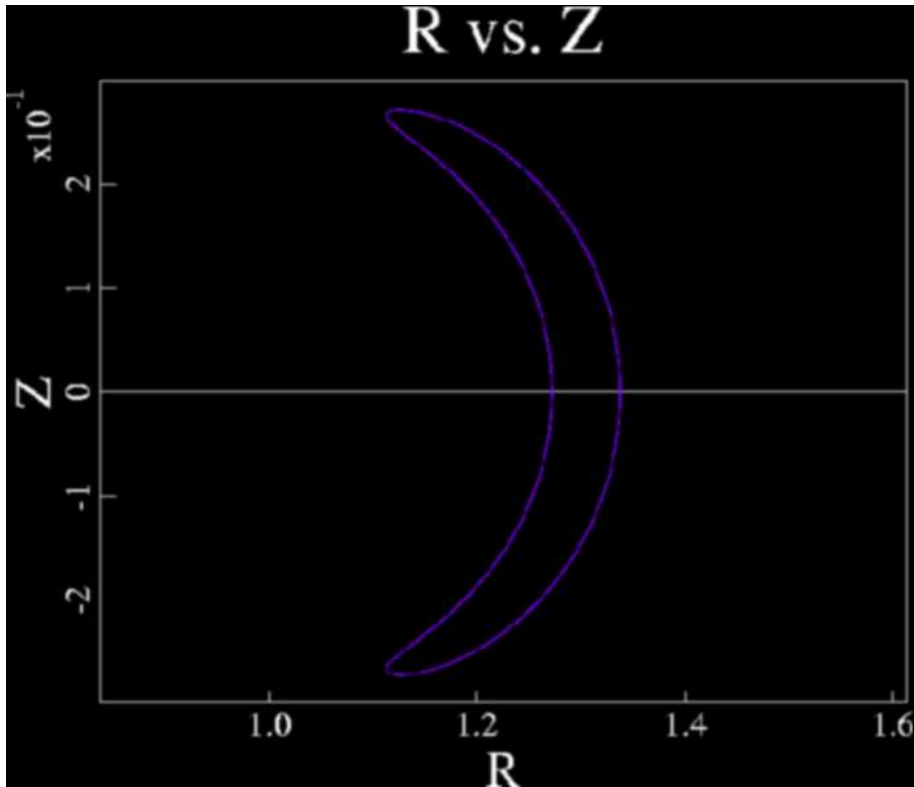
where Δ_j^k is the determinant and k denotes iteration, j denotes FE

- iterative process in **both k and j**
- sorting and parallelization done at the same time
- slowest part!

Examples

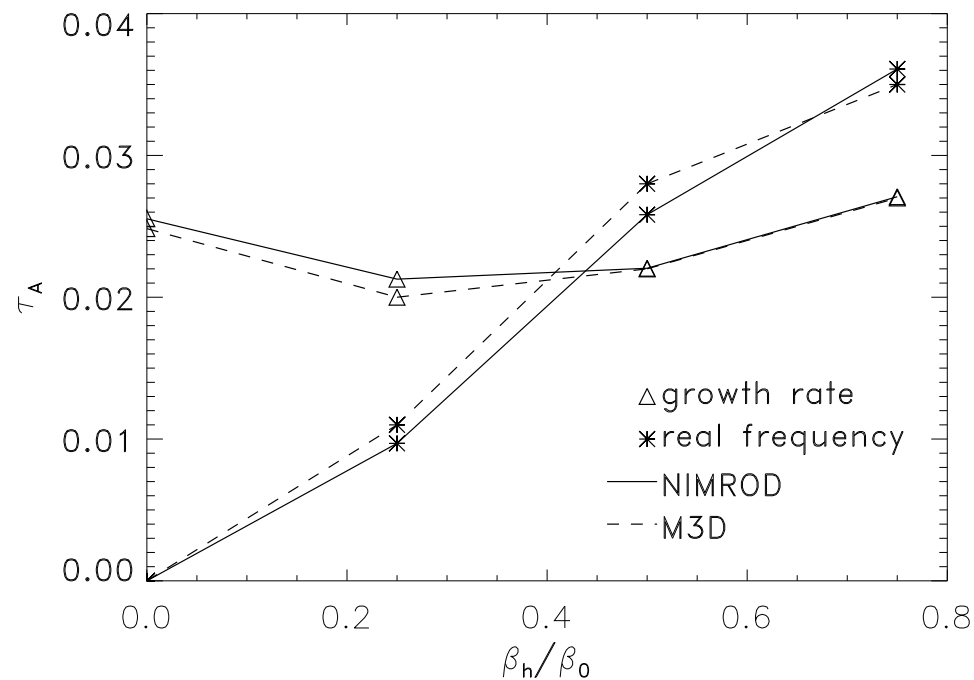


Particle Traces



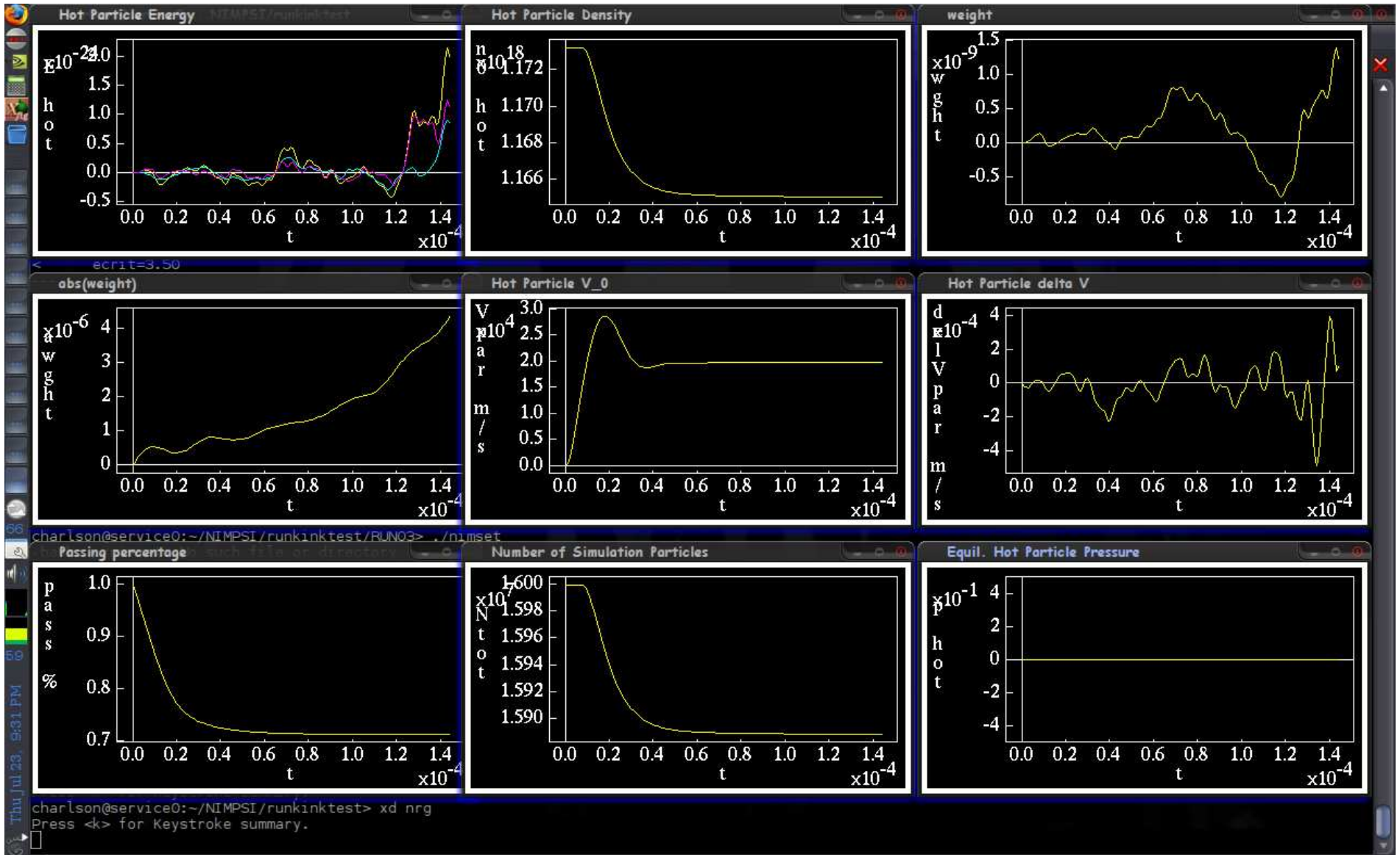
(1, 1) kink benchmark

- monotonic q , $q_0 = .6$, $q_a = 2.5$, $\beta_0 = .08$, shifted circular tokamak $R/a=2.76$
- $dt=1e-7$, $\tau_A = 1.e6$



Diagnostics indicate trajectories are resolved

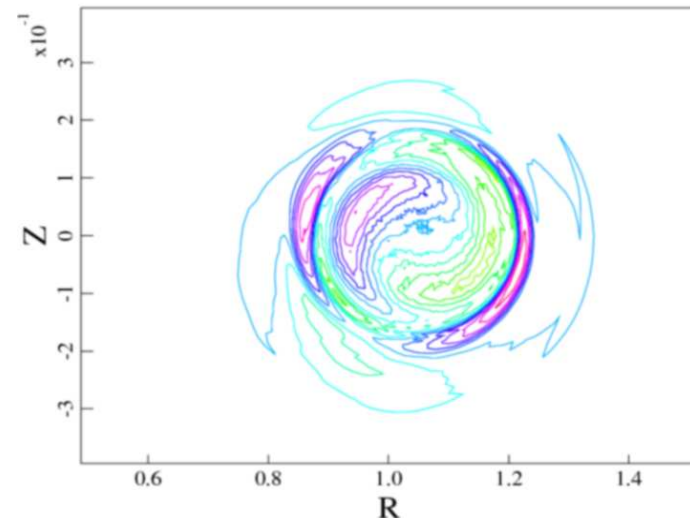
\$>xd nrg



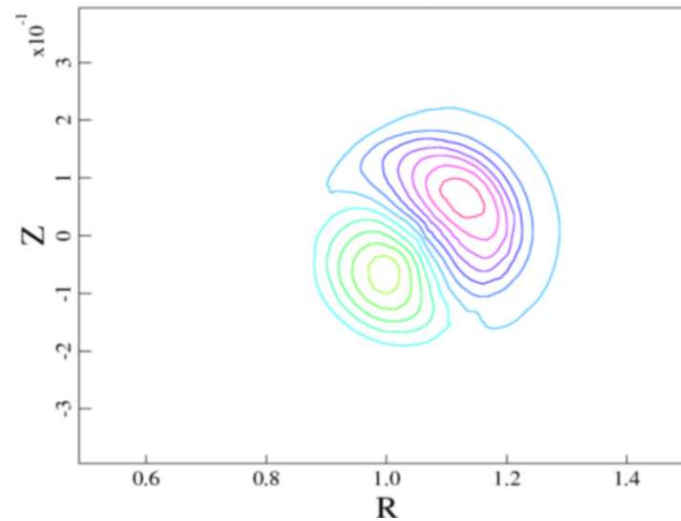
Contours suggest sufficient particle number

\$>xd conph

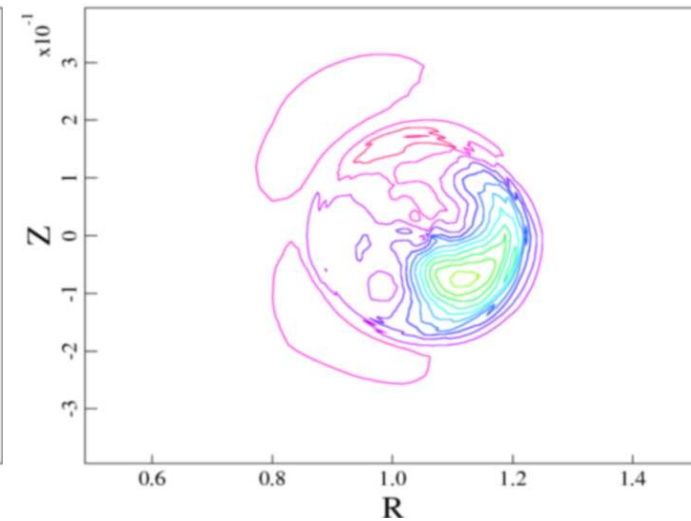
Re VPhi



Re Ph_per

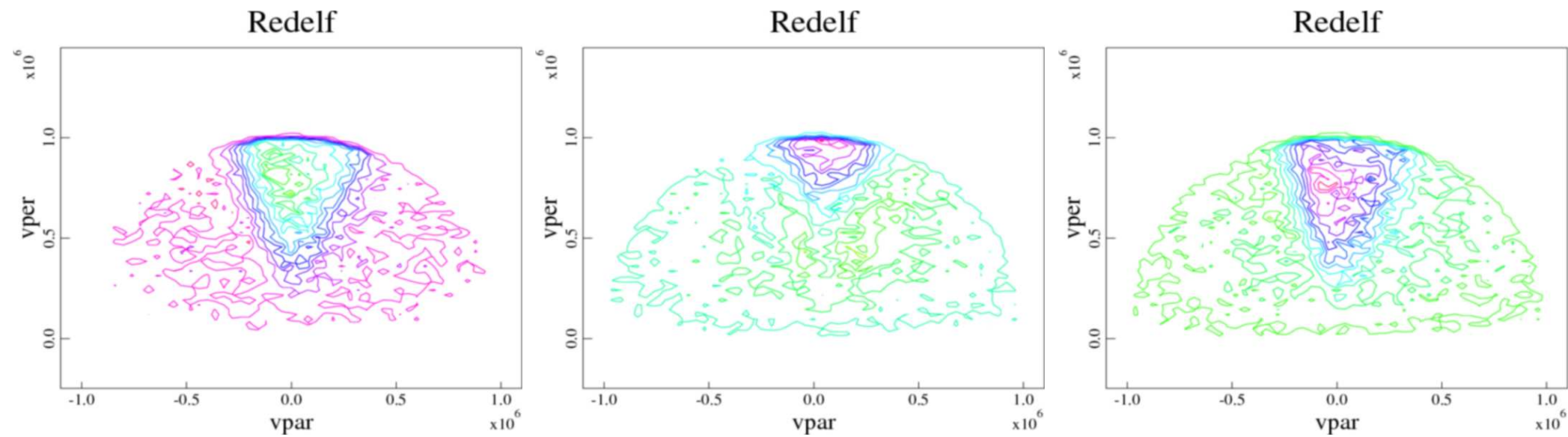


Re Ph_ani



$\delta f(v_{\parallel}, v_{\perp})$ shows activity concentrated in trapped cone

\$>xd phase



at $t = [4.0, 4.5, 5.0] \times 10^{-4} s$

Status of Drift Kinetic Particles

- drift kinetic particles are moderately robust in linear NIMROD
- general parallelization in rblocks
- run on ~ 100 procs with $\sim 10M$'s particles
- up to $\varepsilon_{max} \simeq 250KeV$
- minor bug in restart
- more diagnostics in the works



Future Directions

- better field gather routine **E, B**
- domain decomposition with `nlayer`
- high order finite elements
- improved deposition
- improved parallelization - domain cloning, asynchronous
- utilize sorted particle list
- nonlinear - center equations
- try type of arrays data structure
- full f option
- TAE's and EPM's



Deposition of δp_h onto Finite Element grid

- recall pressure moment at a position \mathbf{x} is

$$\delta p(\mathbf{x}) = \int m(v - V_h)^2 \delta f(\mathbf{x}, \mathbf{v}) d^3v$$

$$\delta p_n(R, Z) = \sum_{i=1}^N m(v_i - V_h)^2 g_0 w_i \delta(R - R_i) \delta(Z - Z_i) e^{-in\phi_i}$$

where sum is over the particles, m mass of the particle, $g_0 w_i$ is the perturbed phase density, and V_h is the hot flow velocity, deposition is done in Fourier space^a

- express lhs $\delta p_n(R, Z)$ in the finite element basis,

$$\delta p_n(R, Z) = \sum_k \delta p_n^k N^k(\eta, \xi)$$

where sum k is over the basis functions and (η, ξ) are functions of (R, Z)

^afor large n runs, a conventional ϕ deposition with a **FFT** is performed

- project onto the finite element space by casting in **weak form**:

$$\int N^l \sum_k \delta p_n^k N^k d^3x = \int N^l \sum_{i=1}^N m(v_i - V_h)^2 g_0 w_i \delta(R - R_i) \delta(Z - Z_i) e^{-in\phi_i} d^3x$$

$$\underline{\mathbf{M}} \delta p_n^k = \sum_{l \in k} \sum_{i=1}^N m(v_i - V_h)^2 g_0 w_i e^{-in\phi_i} N^l(\eta_i, \xi_i)$$

\mathbf{M} is the finite element mass matrix

- **Petrov-Galerkin method**
- **gather** is done using the same shape functions

$$A(\mathbf{x}_i) = \sum_l A^l N^l(\eta_i, \xi_i)$$

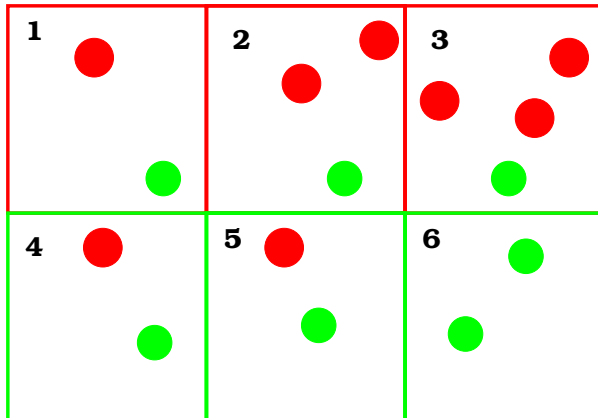
for some field quantity A

- should use shape function $S(R - R_i)$ but requires deposition in real space
 - need ghost cells or global grid with `MPI_ALLGATHER`
 - consistent with **gather**???

Particle Sorting

- Sorting is important because:
 - essential to domain decomposition
 - cache thrashing is minimized
- use “bucket” sort
- particles are sorted on global FE grid
- particles outside of the processor sub-domain (rblock) are passed to their appropriate processor

Bucket Sort

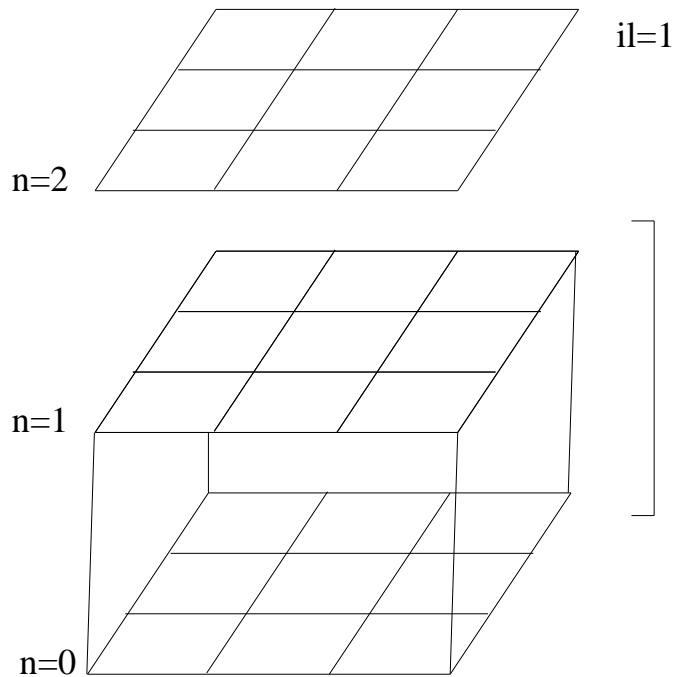


cell number	1	2	3	4	5	6
red count	1	2	3	1	1	0
green count	1	1	1	1	1	2
total count	2	3	4	2	2	2
local displacement	2	5	9	2	4	6

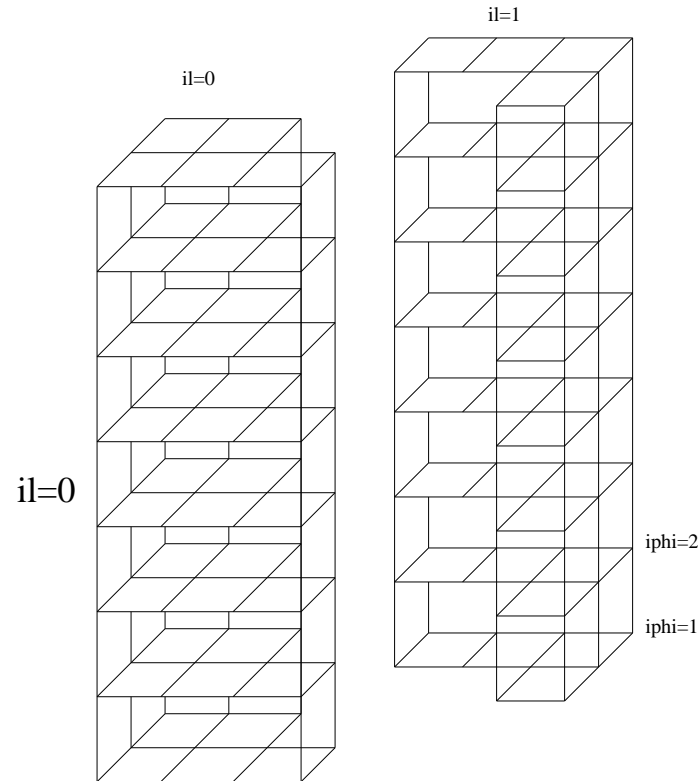
- number of particles in each element is tabulated (**red count**, **green count**)
- **total count** gives the number of particles in each element
- displacement array is calculated from the **total count** - bucket size is determined
- pass particles outside of subdomain (**red particle** to **green cells** and **green particles** to **red cells**)
- using the displacement array, the sorted particle list is filled

Domain Decomposition with $n_{\text{layer}} > 1$

Normal Decomp

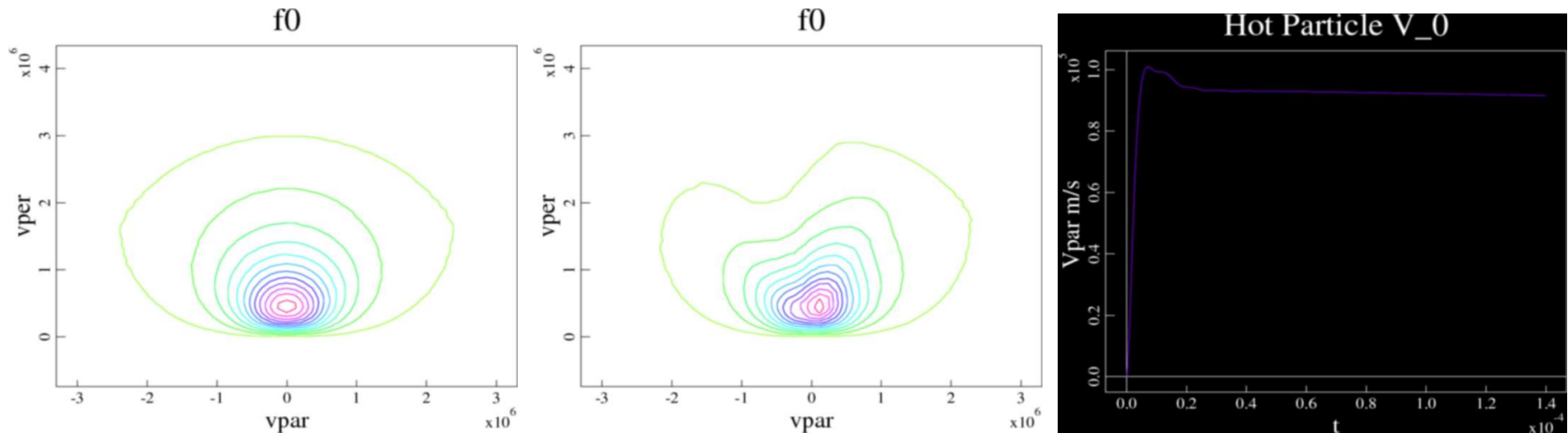


Configuration-space Decomp.



- should not be too bad
- works with NIMROD's domain decomposition
- forces deposition to be done in real space
- particles will reside in n_{layer} subdomain of rblocks

Effect of ignoring $g\rho_{\parallel}$



- tricky to take into account due to mixing of \mathbf{x} and \mathbf{v}
 - inclusion results in anisotropy
 - equilibrium balance is no longer with scalar pressure
- mostly a sampling issue
- maybe a consistency issue
- $V_{hot} \neq 0!$ in few transit times
- f_0 **NOT** isotropic

TAE's

- D. Spong for ITPA Meeting
- toroidicity induce Alfvén eigenmodes (TAE) are marginally stable global modes that exist in the toroidicity induced gap of the Alfvén continuum
- energetic particles can destabilize
- have equilibrium specifications
 - shifted circular equilibrium
 - zero β , aspect ratio .1
 - $q(\rho) \simeq 1.1 + 0.8\rho^{2.4}$
 - energetic particles are Maxwellian $T_h \simeq 10\text{KeV}$
- may impact heating methods and confinement of α 's
- TAE's also of interest to RFP's
- possible evidence of TAE's in extensions of benchmark case

