

The effect of poloidal flow on locked islands

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Plasma flow is observed to “heal” magnetic islands.

Recent work at the Large Helical Device (LHD) indicates that plasma flow plays a primary role in “healing” vacuum magnetic islands in stellarators. An analytic theory to explain this phenomenon has been proposed. ¹

- The model to explain this is reminiscent of ‘mode locking/unlocking’ physics in tokamaks.
- Governed by a torque balance between viscous and electromagnetic torques:

$$D_A \sin(2\Delta\phi) = D_\Omega$$

where:

$$D_A = \omega_A^2 \frac{m_o^2 k_v^2}{k_f c_r} \frac{(-\rho_0 \Delta'_0)}{2m_0} \frac{(t'_0)^2 (w^V)^4}{512}$$

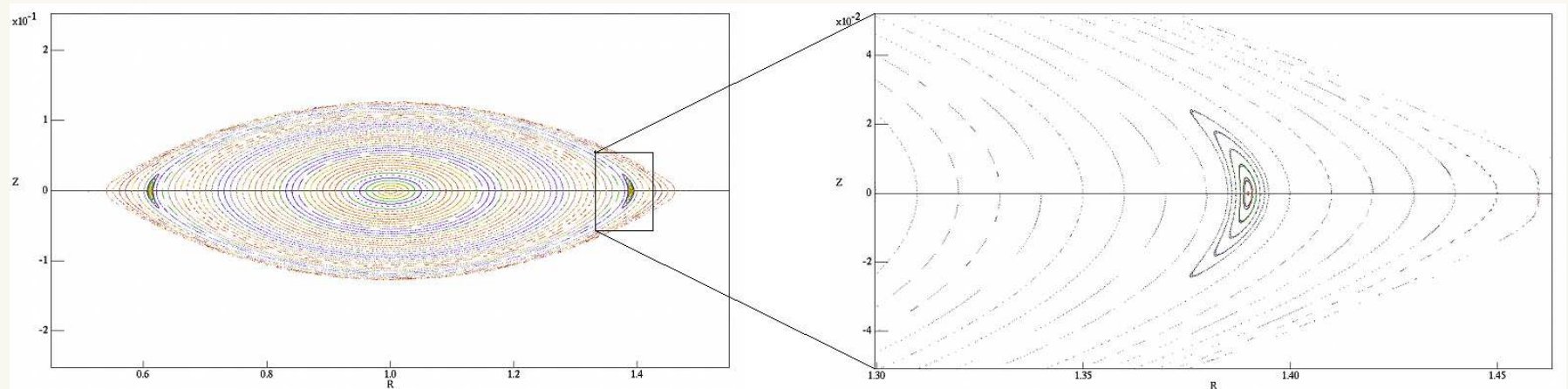
$$D_\Omega = -2\nu_\perp \frac{\rho_0}{\delta_r} \Omega_0^\alpha$$

¹Hegna, UW-CPTC Report 11-7, 2011.

“Healing” of magnetic islands is investigated with NIMROD.

- A $m=2, n=2$ straight stellarator vacuum field is created using tools developed for previous work.
- This field is perturbed with $m=2, n=5$ symmetry-spoiling harmonics which create islands in the vacuum field. Magnetic potential for this situation is given by:

$$\phi = B_0 \left[R\zeta + \epsilon_{22} \frac{2R}{2} I_2 \left(\frac{2r}{R} \right) \sin(2\theta - 2\zeta) + \epsilon_{25} \frac{2R}{5} I_2 \left(\frac{5r}{R} \right) \sin(2\theta - 5\zeta) \right]$$



Flow is achieved by using an ad-hoc injection of poloidal momentum.

A source of poloidal momentum is introduced, which is aligned with flux surface projections on the poloidal plane. This alignment is accomplished by:

$$\mathbf{f} \cdot \nabla\psi = 0$$

where

$$\mathbf{f} = f^R \overline{\mathbf{e}}_R + f^Z \overline{\mathbf{e}}_Z$$

is the injected poloidal momentum source.

Specifically,

$$\sqrt{(f^R)^2 + (f^Z)^2} = v_{\text{cent}}$$
$$f^R \frac{\partial\psi}{\partial R} + f^Z \frac{\partial\psi}{\partial Z} = 0$$

where v_{cent} is a parameter specified in `nimrod.in` and represents the magnitude of the injected poloidal momentum source.

The poloidal momentum source does indeed produce velocity vectors aligned with the flux surface projection in the poloidal plane.

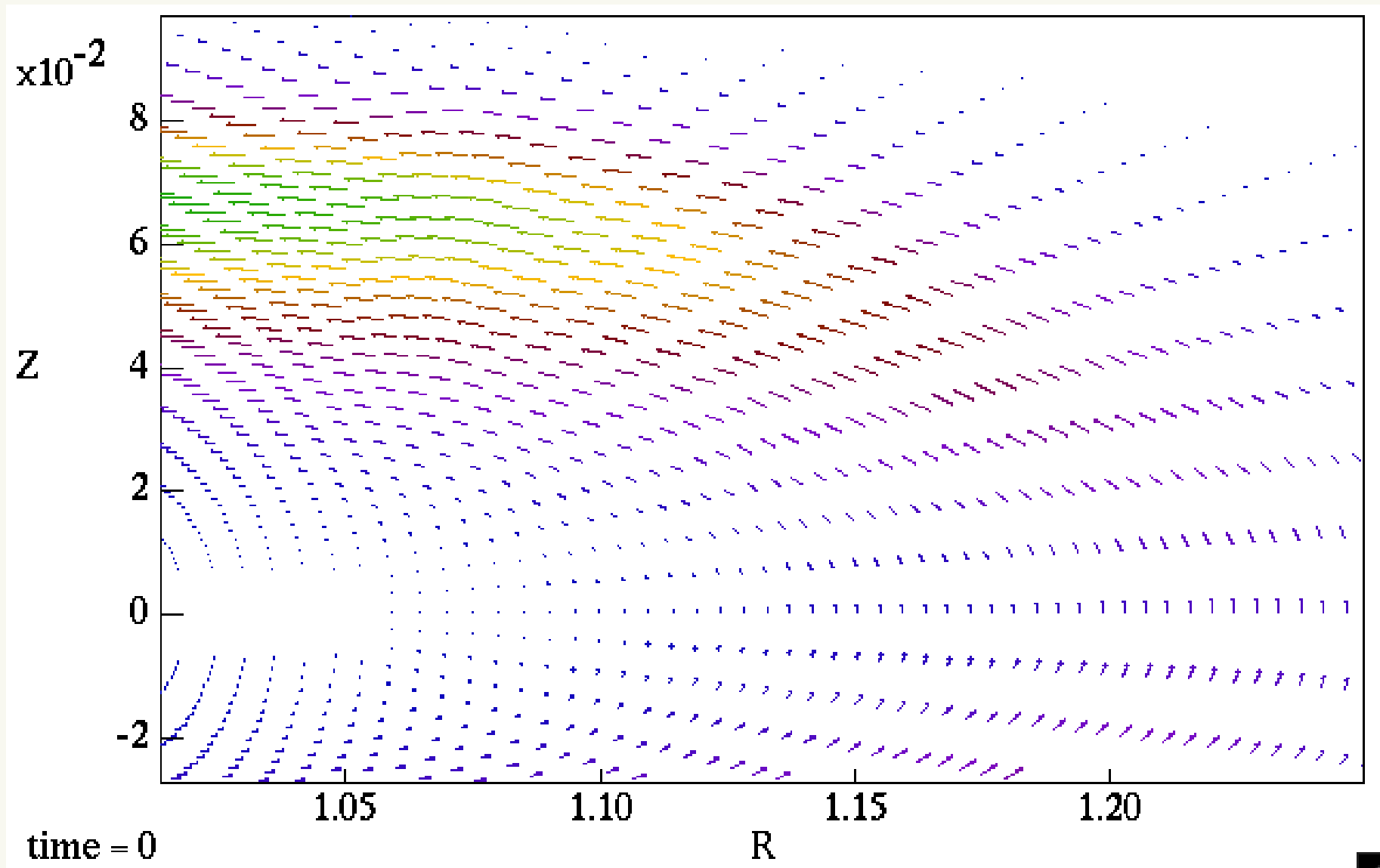


Figure 1: Velocity vectors at $t=0.2$ ms.

Velocity profiles for the $(m=2, n=2, \epsilon=0.775)$ case, slow ramping of injected poloidal momentum.

The island for this case is between $R=1.26$ and $R=1.29$.

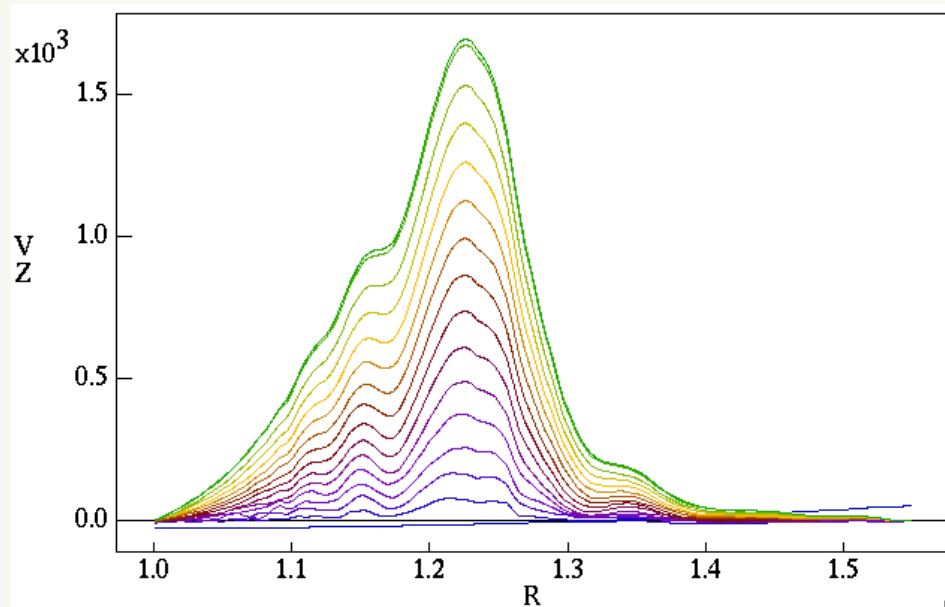


Figure 2: V_Z profiles at various times, $\zeta = 0$.

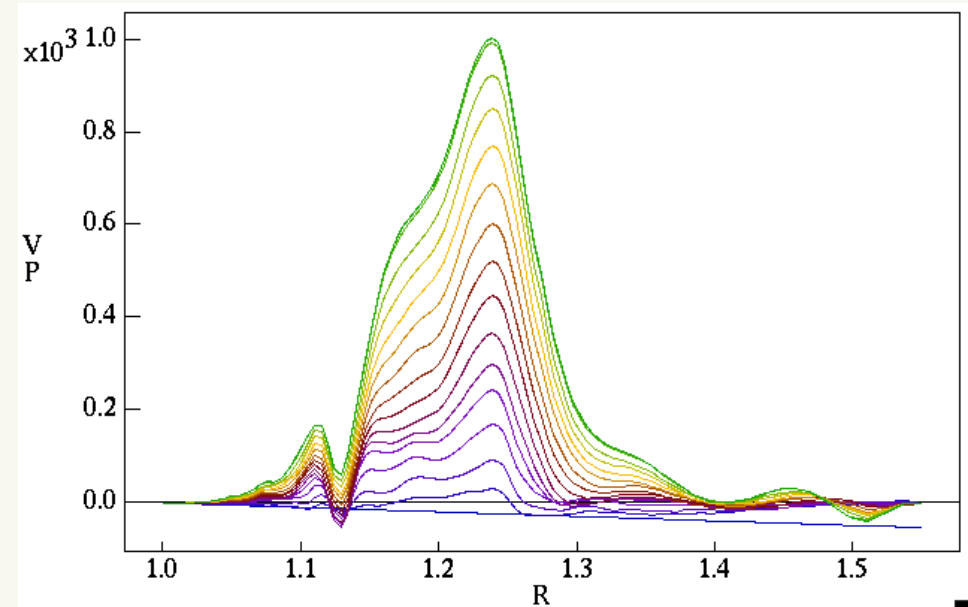


Figure 3: V_ϕ profiles at various times, $\zeta = 0$.

Results for the $(m=2, n=2, \epsilon=0.775)$ case, slow ramping of injected poloidal momentum.

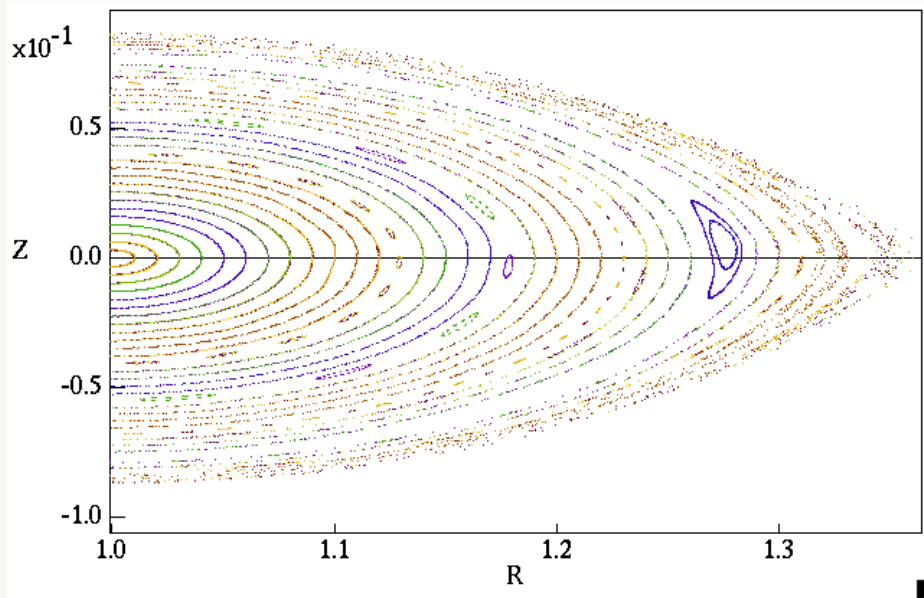


Figure 4: Poincare plot at $t=0.061\text{ms}$, $\zeta = 0$.

Simulation: $\Delta\phi = 1.21^\circ$.

Hegna prediction: $\Delta\phi = 0.4^\circ$.

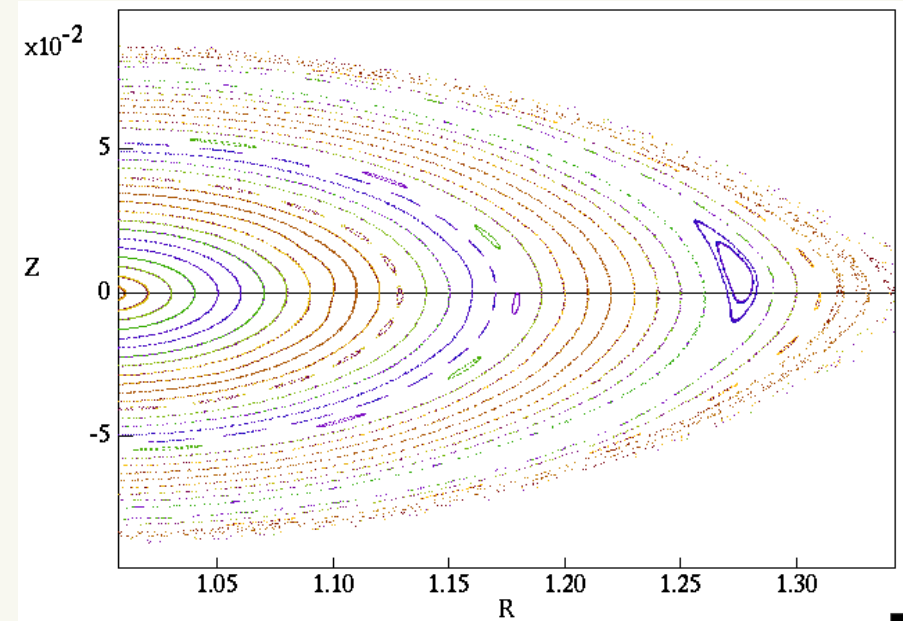


Figure 5: Poincare plot at $t=0.152\text{ms}$, $\zeta = 0$.

Simulation: $\Delta\phi = 1.61^\circ$.

Hegna prediction: $\Delta\phi = 1.0^\circ$.

Velocity profiles for the ($m=2, n=2, \epsilon=0.775$) case, fast ramping of injected poloidal momentum.

The island for this case is between $R=1.26$ and $R=1.29$.

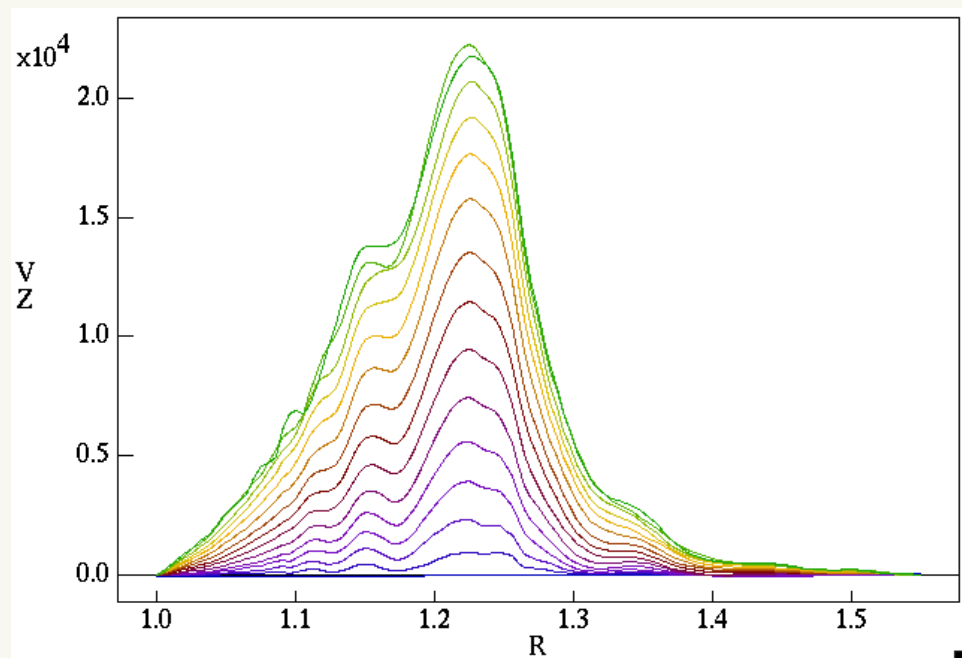


Figure 6: V_Z profiles at various times, $\zeta = 0$.

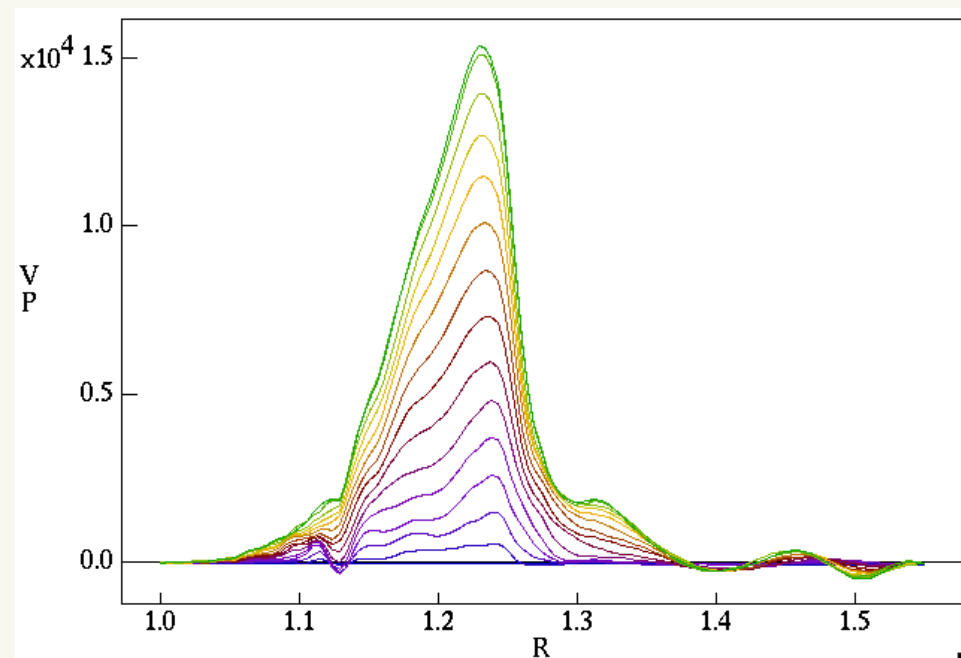


Figure 7: V_ϕ profiles at various times, $\zeta = 0$.

Results for the $(m=2, n=2, \epsilon=0.775)$ case, fast ramping of injected poloidal momentum.

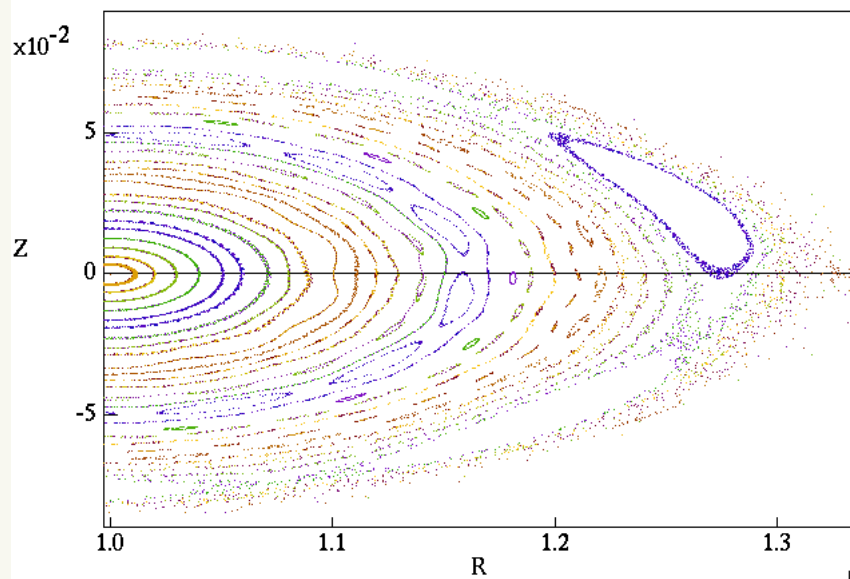


Figure 8: Poincare plot at $t=0.089\text{ms}$, $\zeta = 0$.

Simulation: $\Delta\phi = 6.34^\circ$.

Hegna prediction: $\Delta\phi = 6.3^\circ$.

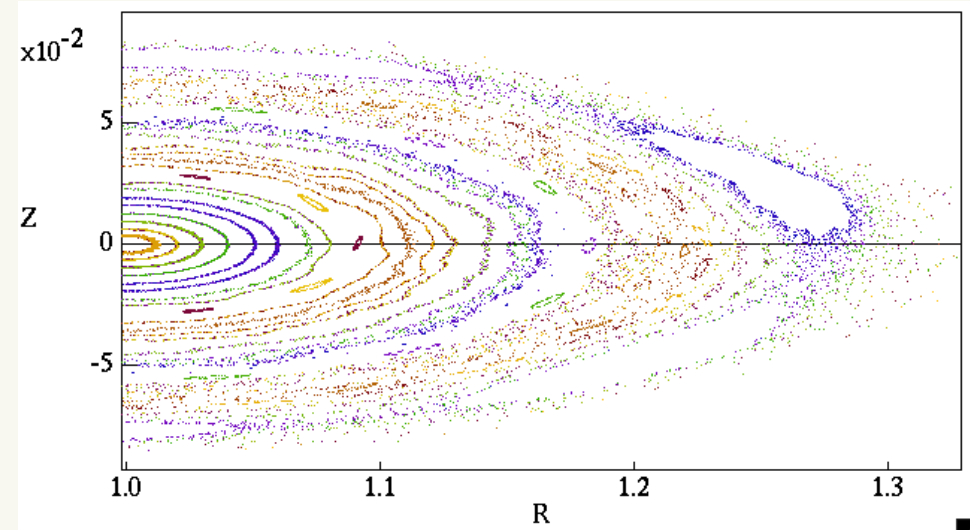


Figure 9: Poincare plot at $t=0.214\text{ms}$, $\zeta = 0$.

Simulation: $\Delta\phi = 6.46^\circ$.

Hegna prediction: $\Delta\phi = 13.8^\circ$.

Results for the $(m=2, n=2, \epsilon=0.80)$ case, slow ramping of injected poloidal momentum.

The island for this case is between $R=1.22$ and $R=1.24$.

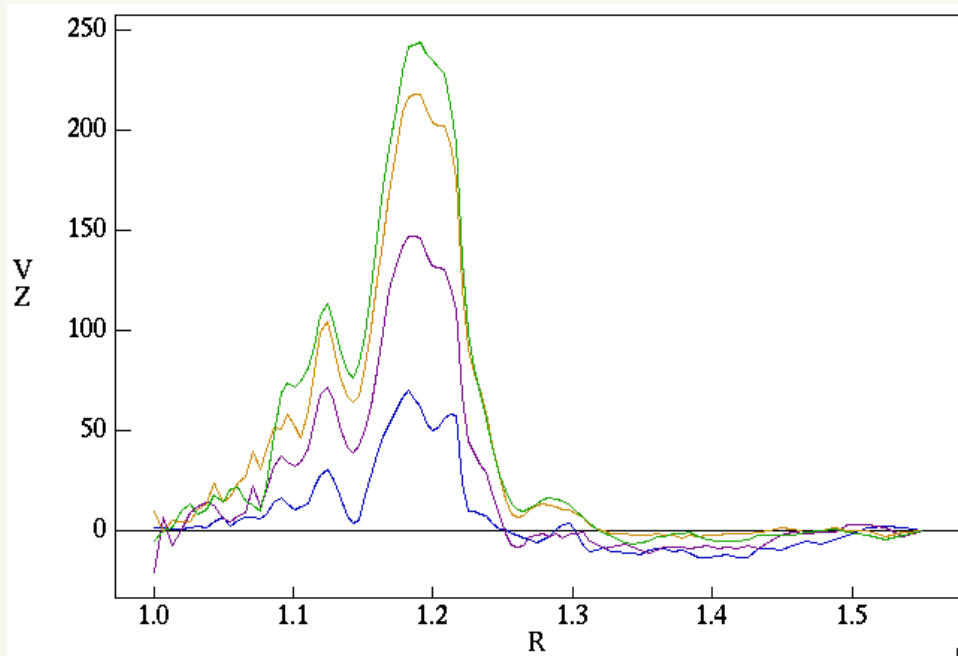


Figure 10: V_Z profiles at various times, $\zeta = 0$.

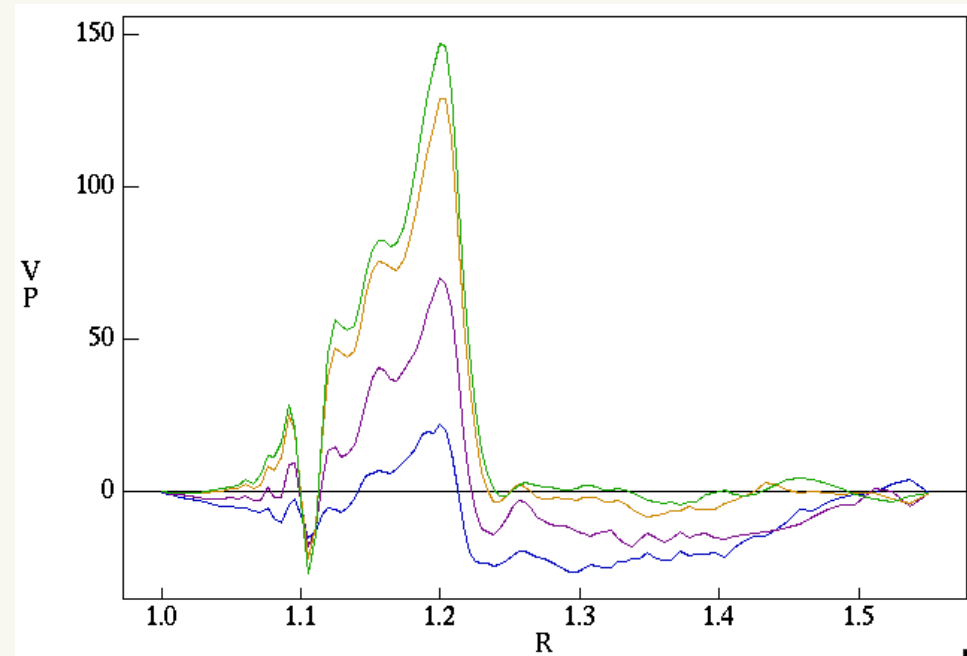


Figure 11: V_ϕ profiles at various times, $\zeta = 0$.

Results for the $(m=2, n=2, \epsilon=0.80)$ case, slow ramping of injected poloidal momentum.

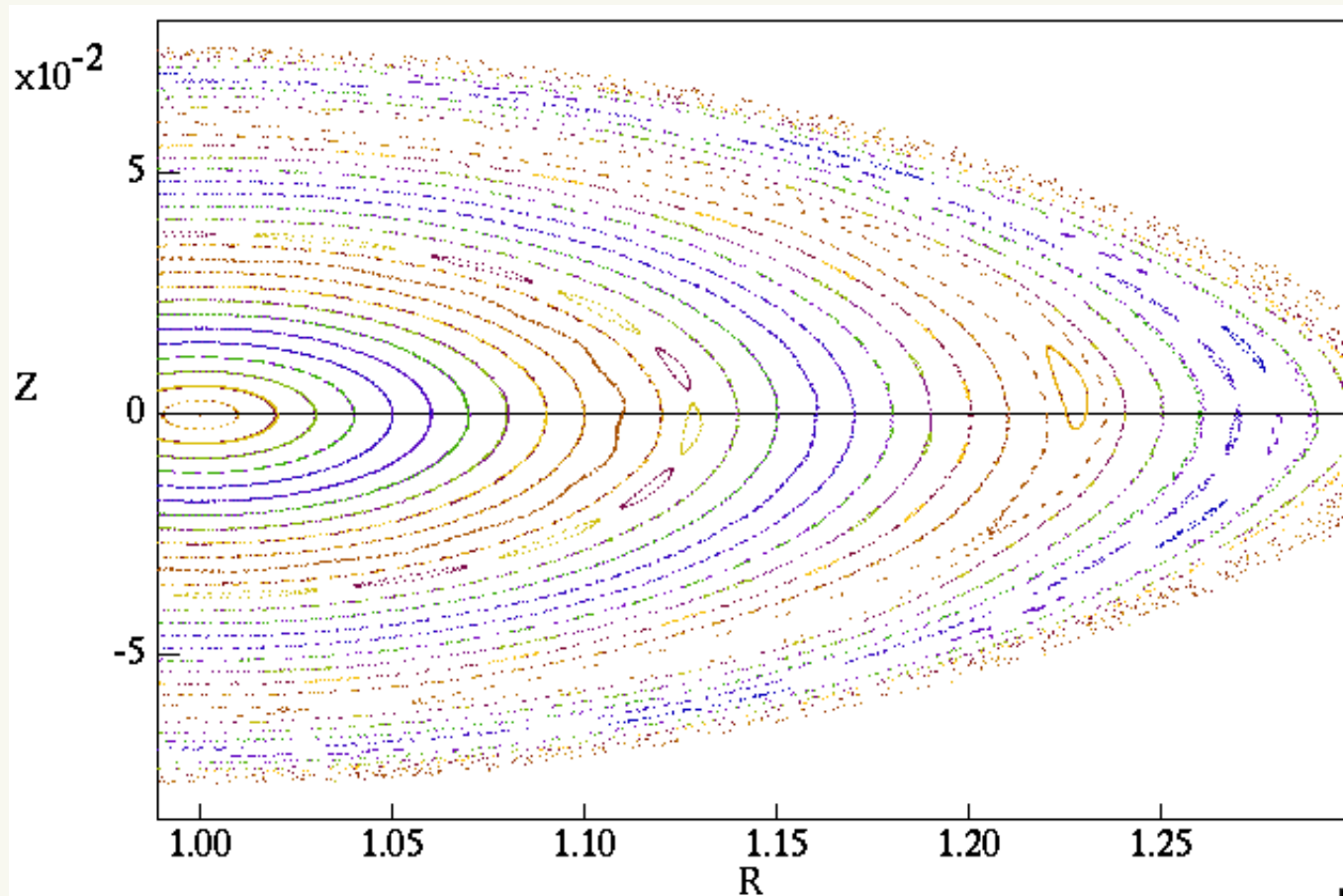


Figure 12: Poincaré plot at $t=0.018$ ms.

Simulation: $\Delta\phi = 1.55^\circ$.

Hegna prediction: $\Delta\phi = 1.1^\circ$.

Results for the $(m=2, n=2, \epsilon=0.80)$ case, fast ramping of injected poloidal momentum.

The island for this case is between $R=1.22$ and $R=1.24$.

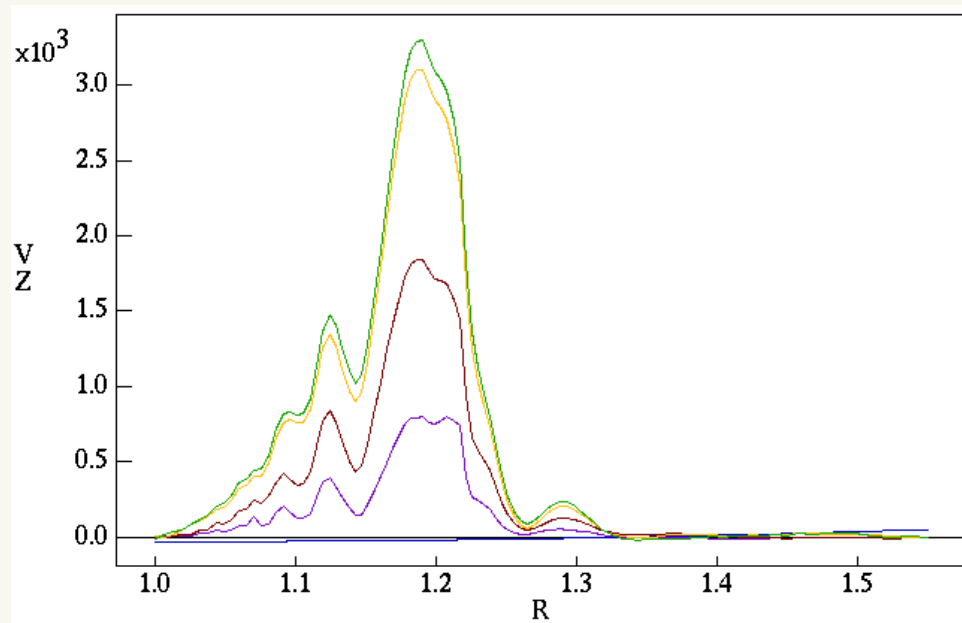


Figure 13: V_Z profiles at various times, $\zeta = 0$.

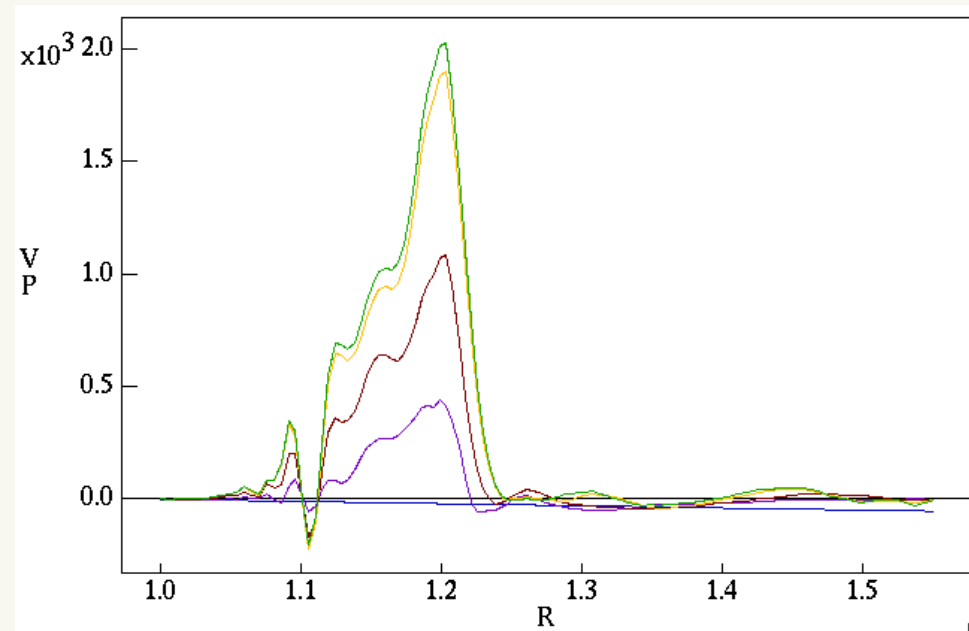


Figure 14: V_ϕ profiles at various times, $\zeta = 0$.

Results for the ($m=2, n=2, \epsilon=0.80$) case, fast ramping of injected poloidal momentum.

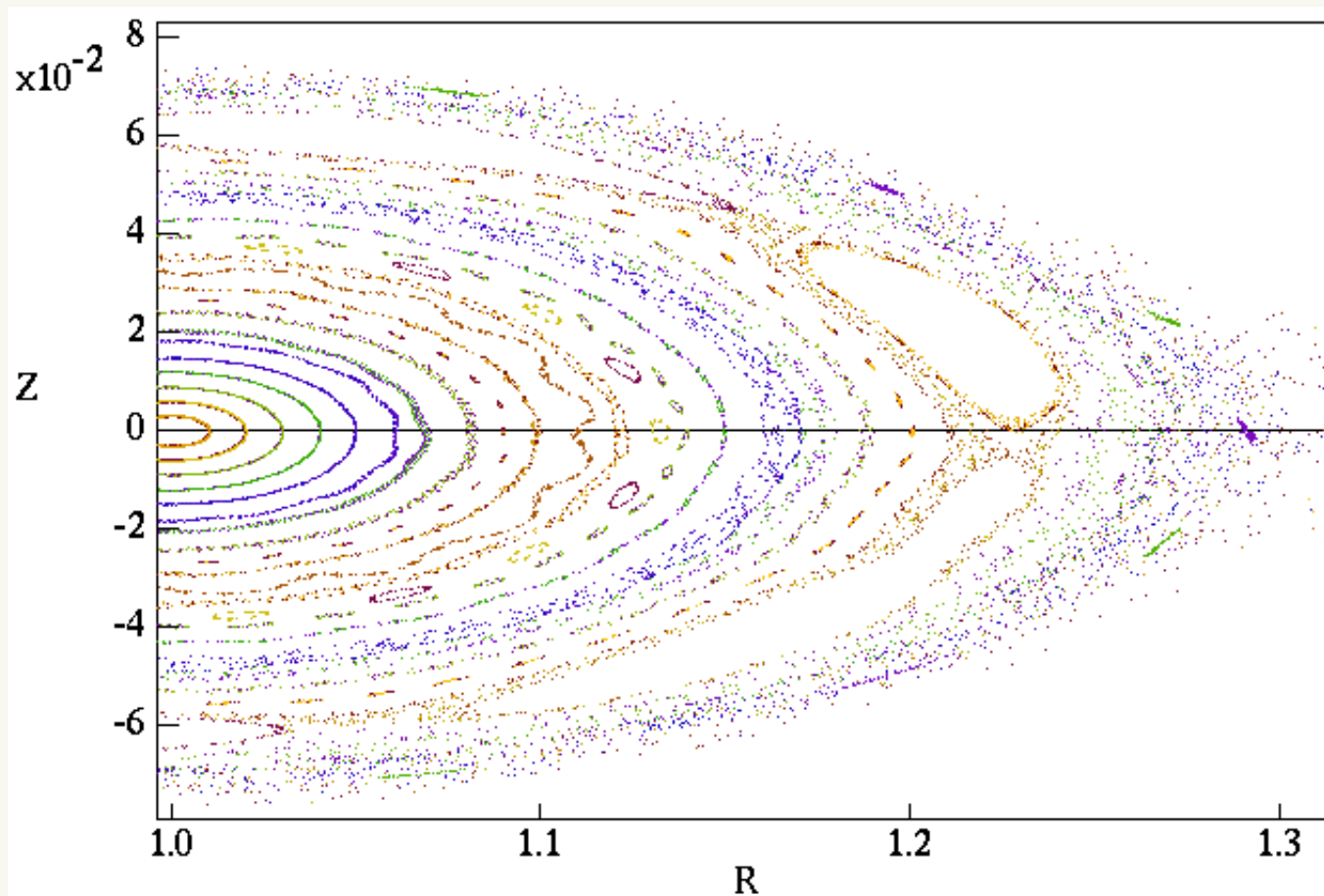


Figure 15: Poincaré plot at $t=0.025$ ms.

Simulation: $\Delta\phi = 6.16^\circ$.

Hegna prediction: $\Delta\phi = 17.2^\circ$.

Summary of results.

- It appears that the simulation and analytic theory are close for small flows. For larger flows, the agreement deteriorates.
- Possible sources of disagreement:
 - Evaluation of geometrical parameters - analytic work is for a cylinder. Flux surfaces in the simulation are not circular.
 - Evaluation of velocities and radii - velocity is not constant on a flux surface projection.
 - Much flux surface destruction is seen in the very high flow cases. Does the model still hold?
 - In the very high flow cases, the flow speeds are exceeding the sound speed in the simulation.

Future work:

- Run cases with the islands placed even further toward the core.
- Increase the background temperature to avoid any shocks.