



NUMERICAL STUDY OF A HIGH BETA DISRUPTION IN DIII-D SHOT 87009 WITH THE NIMROD CODE

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OUTLINE

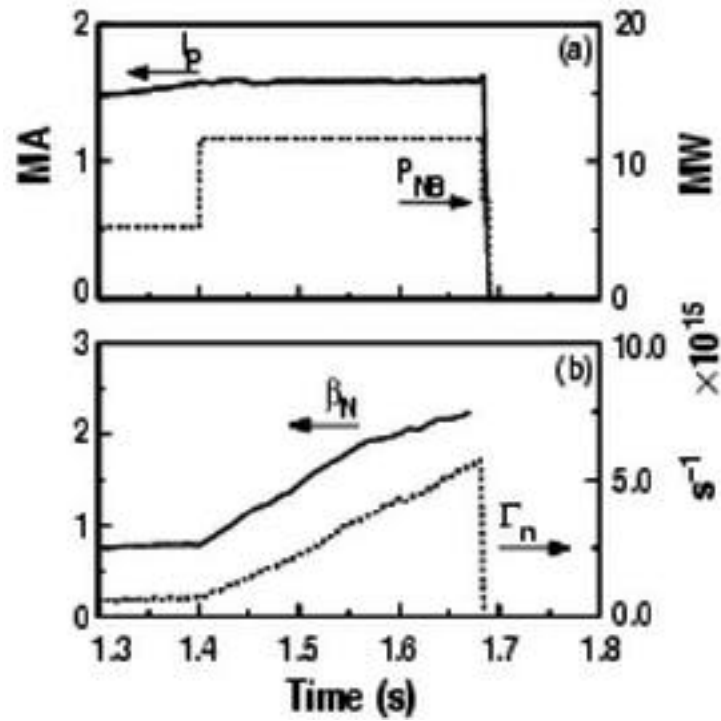
- DIII-D shot 87009
 - Slow heating through critical β_N ($\beta_{Ncrit} \sim 1$) **High- β disruption**
 - Perturbation growth greater than usual $\exp(\gamma t)$
- Theory [Callen, Hegna, Rice, Strait, and Turnbull, Phys. Plasmas 6, 2963 (1999)]
 - Slow heating through ideal MHD stability boundary

$$\beta = \beta_c(1 + \gamma_h t) \quad \gamma(t) = \hat{\gamma}_{MHD} \sqrt{\gamma_h t}, \quad \xi \sim \exp[(t/\tau)^{3/2}], \quad \tau \sim \hat{\gamma}_{MHD}^{-2/3} \gamma_h^{-1/3}$$

- Numerical simulation with **NIMROD**
 - Conducting wall on separatrix **Larger β_{Ncrit} ($\beta_{Ncrit} \sim 4.5$)**
 - Linear stability, comparison with GATO
 - Nonlinear ($n = 0, 1$) simulation at $\beta_N = 5$
 - Slow drive through critical point

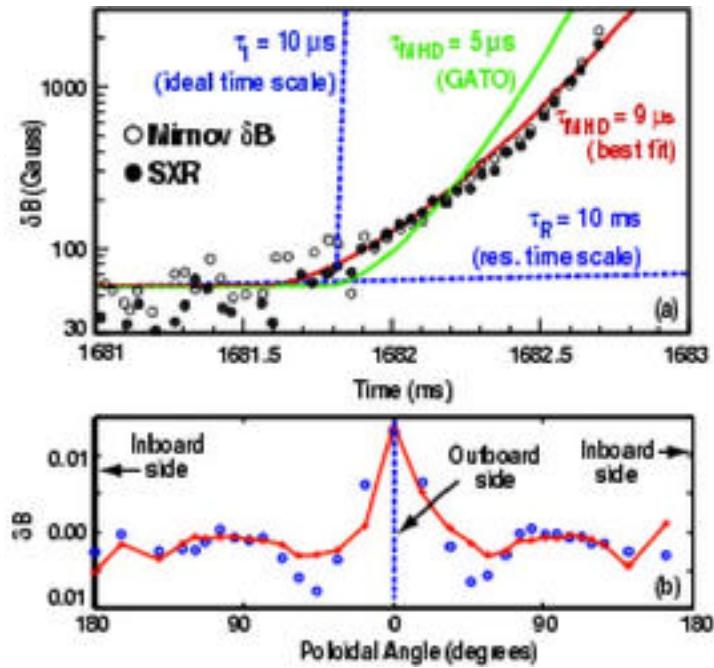
DIII-D SHOT 87009

High- β disruption when heated slowly through critical β_N



DIII-D SHOT 87009

Growth is faster than simple exponential



THEORY

Callen, Hegna, Rice, Strait, and Turnbull, Phys. Plasmas 6, 2963 (1999)

• Heat slowly through critical β : $\beta = \beta_c(1 + \gamma_h t)$

• Ideal MHD: $\omega^2 = -\hat{\gamma}_{MHD}^2 (\beta / \beta_c - 1)$

$$\leftarrow \gamma(t) = \hat{\gamma}_{MHD} \sqrt{\gamma_h t}$$

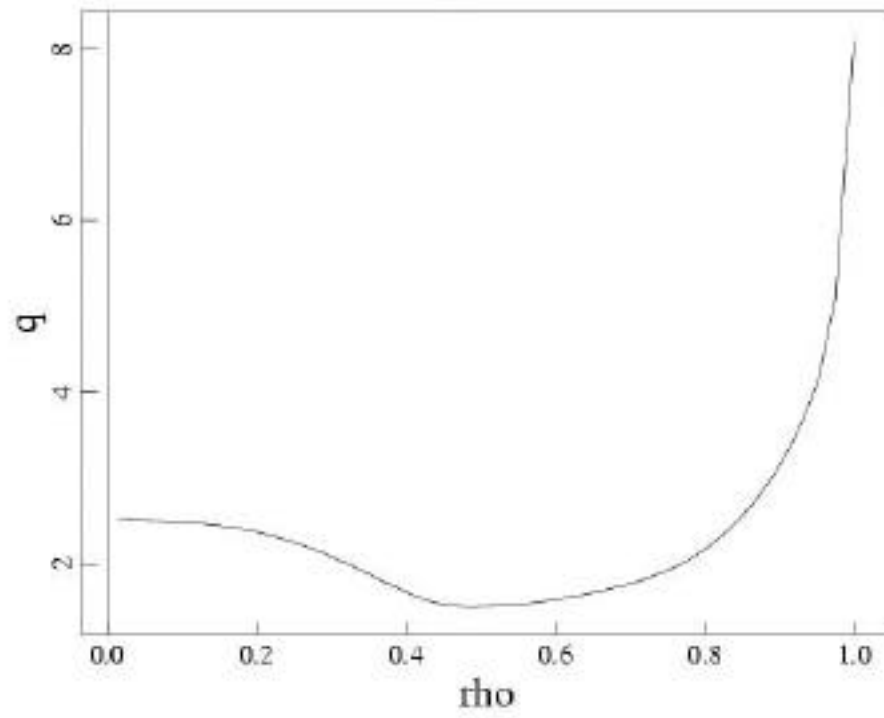
• Perturbation growth:

$$\frac{d\xi}{dt} = \gamma(t)\xi \quad \leftarrow \xi = \xi_0 \exp[(t/\tau)^{3/2}], \quad \tau = (3/2)^{2/3} \hat{\gamma}_{MHD}^{-2/3} \gamma_h^{-1/3}$$

Good agreement with experiment

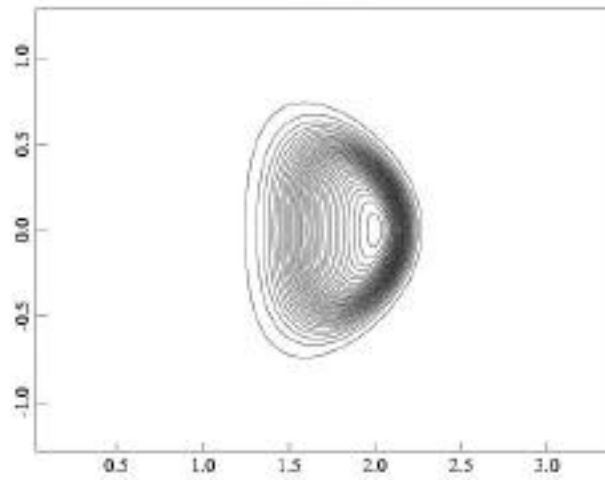
EQUILIBRIUM

q-profile

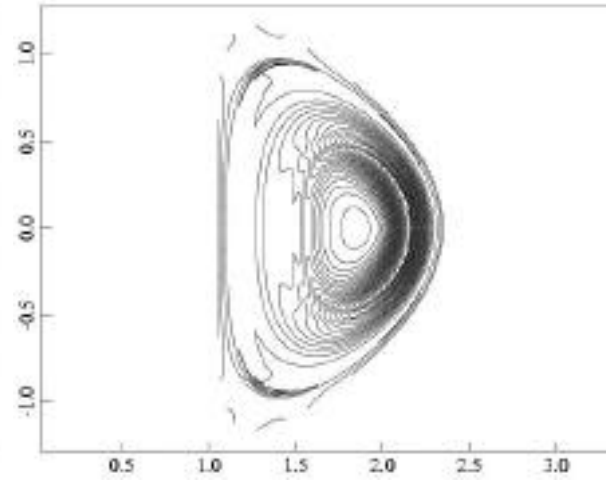


EQUILIBRIUM

Pressure

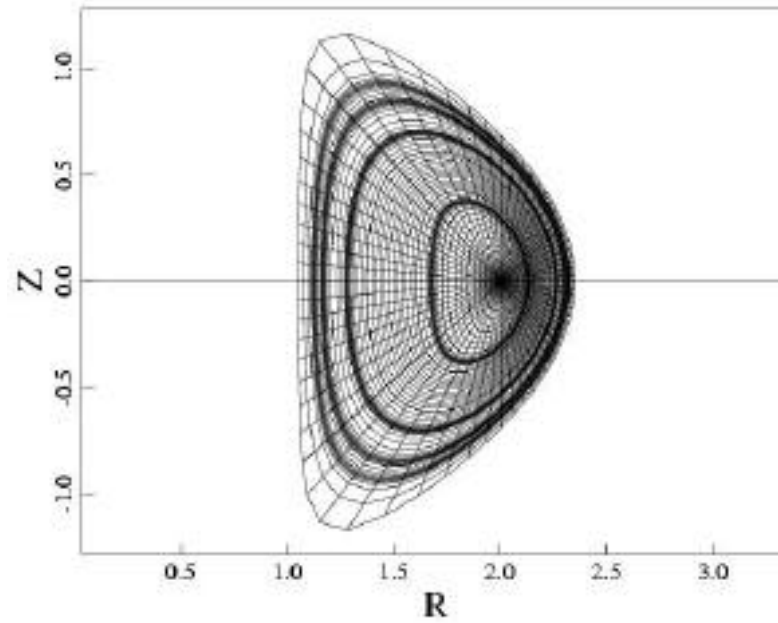


Toroidal Current



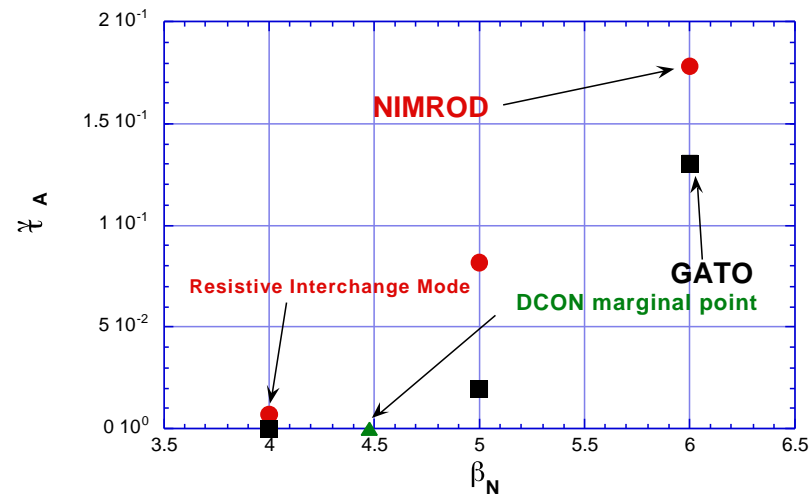
NIMROD GRID

- Packed at rational surfaces
- Bi-cubic finite elements

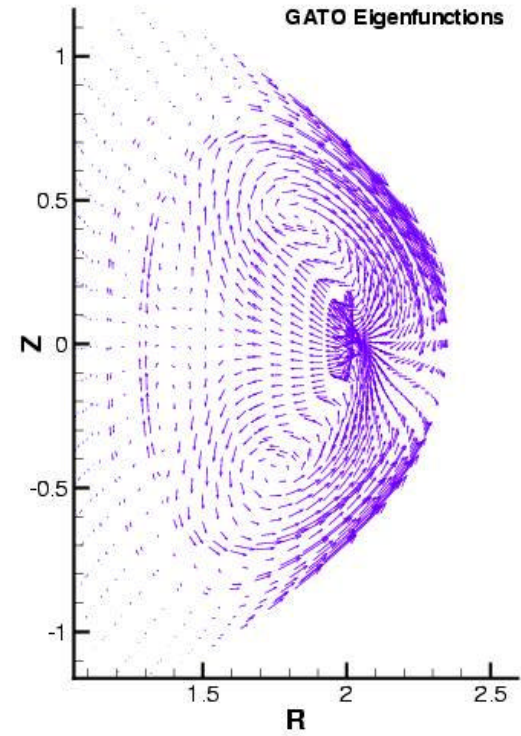
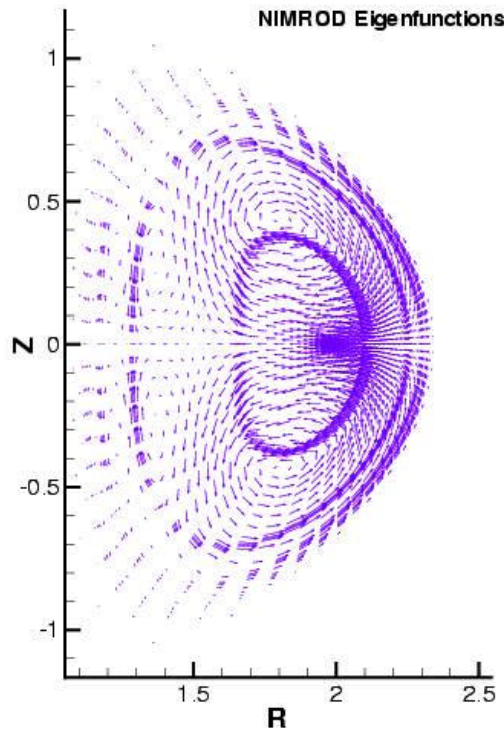


LINEAR STABILITY

- Linear growth rate as a function of β_N from NIMROD and GATO
- NIMROD finds resistive interchange mode below β_{Ncrit}

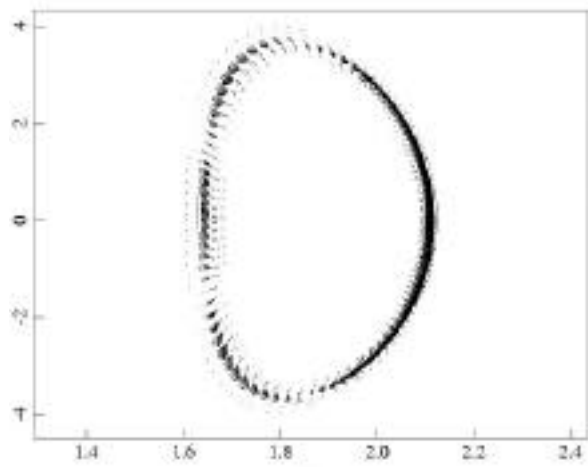


Poloidal Velocity Eigenfunctions

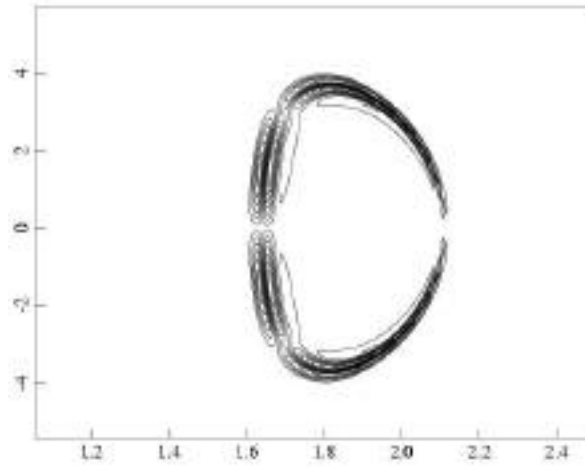


RESISTIVE INTERCHANGE MODE, $\beta_N = 4$

$n = 1$ Poloidal Velocity

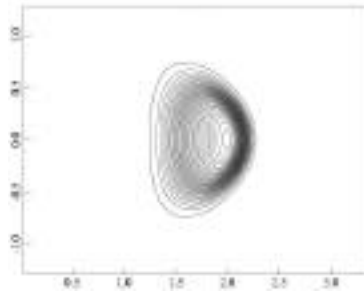


$n = 1$ Pressure



NONLINEAR EVOLUTION, $\beta_N = 5$, PRESSURE

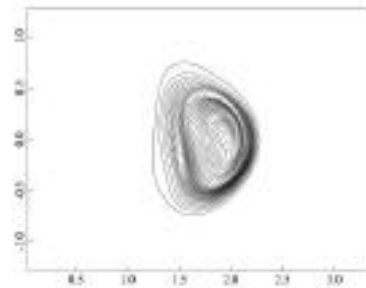
$t = 0$ msec.



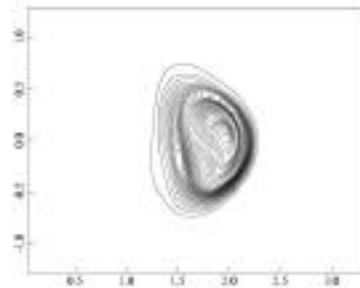
$t = 0.120$ msec



$t = 0.131$ msec

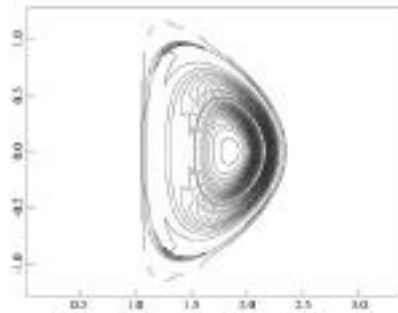


$t = 0.135$ msec.

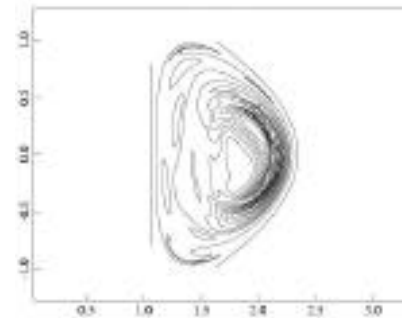


NONLINEAR EVOLUTION, $\beta_N = 5$, TOROIDAL CURRENT

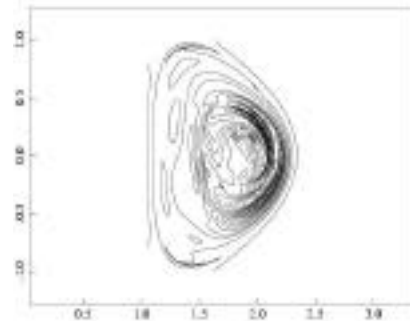
$t = 0$ msec.



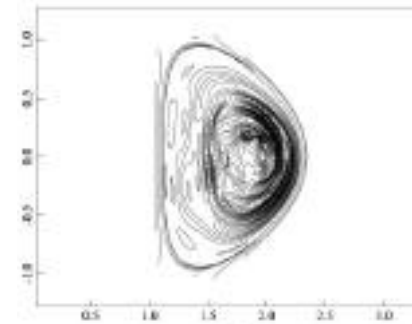
$t = 0.120$ msec



$t = 0.131$ msec

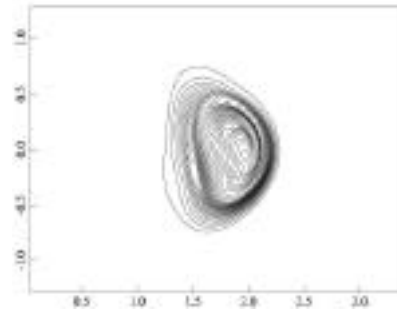


$t = 0.135$ msec.

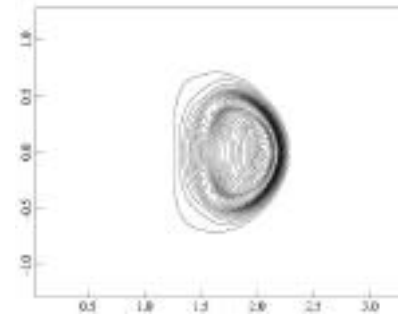


NONLINEAR PRESSURE VS. TOROIDAL ANGLE

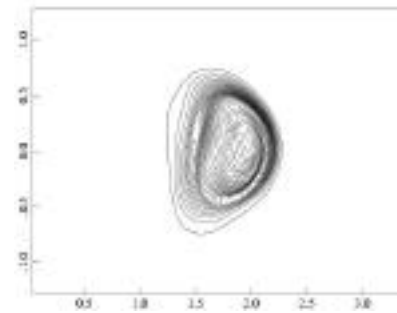
$\phi = 0$



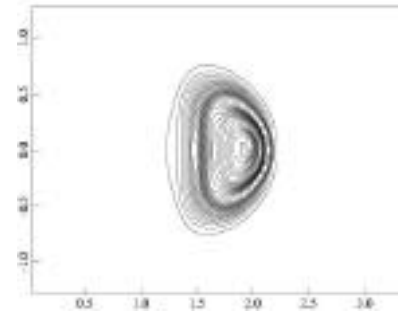
$\phi = \pi/2$



$\phi = \pi$



$\phi = 3\pi/2$



NONLINEAR SIMULATION WITH HEATING

- Initial condition: equilibrium below ideal marginal β_N

- NIMROD $\implies \beta_{Nc} \sim 4.85$ DCON $\implies \beta_{Nc} \sim 4.45$

- Impose heating source proportional to equilibrium pressure profile

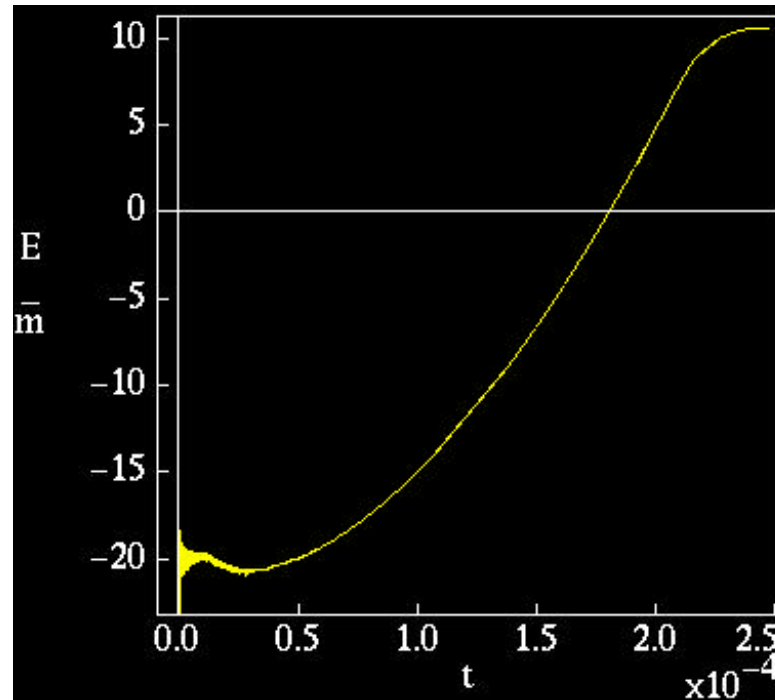
$$\frac{\partial P}{\partial t} = \dots + \gamma_H P_{eq} \implies \beta_N = \beta_{Nc} (1 + \gamma_H t)$$

- Follow evolution through heating, destabilization, and saturation

ENERGY IN UNSTABLE MODE

Magnetic Energy in $n = 1$ mode versus time

$$S = 10^6, \text{ Pr} = 200, \gamma_H = 10^3 \text{ sec}^{-1}$$

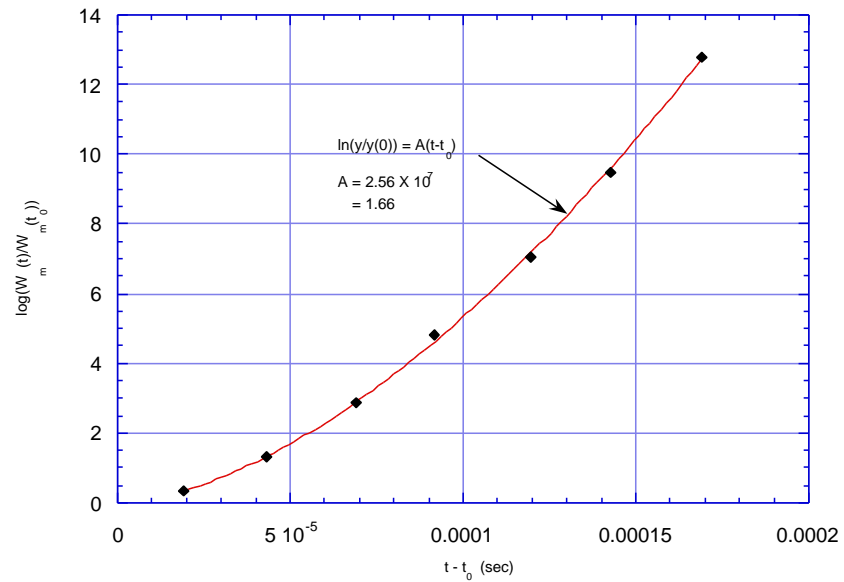


ENERGY GROWS "SUPER-EXPONENTIALLY"

DIII-D Shot #87009 $\log W_m$ versus time

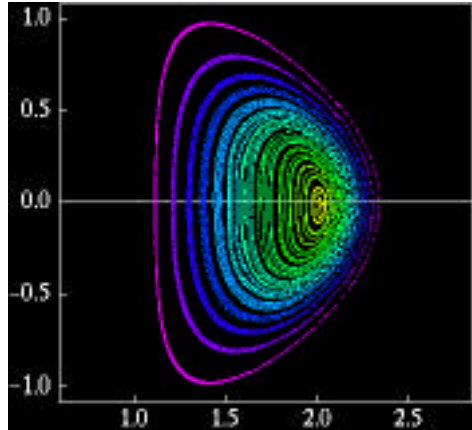
$$\gamma_{\text{HEAT}} = 10^3 \text{ sec}^{-1}, \gamma_{\text{MHD}} = 6.57 \times 10^5 \text{ sec}^{-1}, \beta_N(0) = 4.7$$

$S = 10^6$, $Pr = 200$, Grid: 128 X 64 X 2, Bi-cubic elements

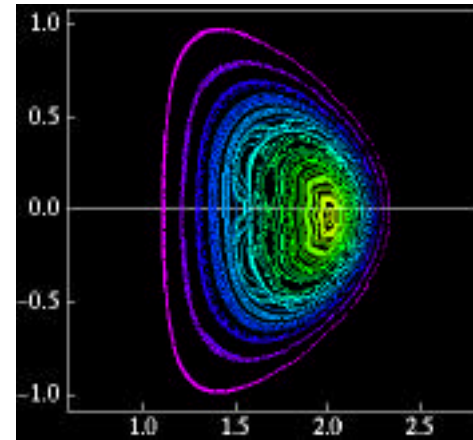


EVOLUTION OF FIELD LINES

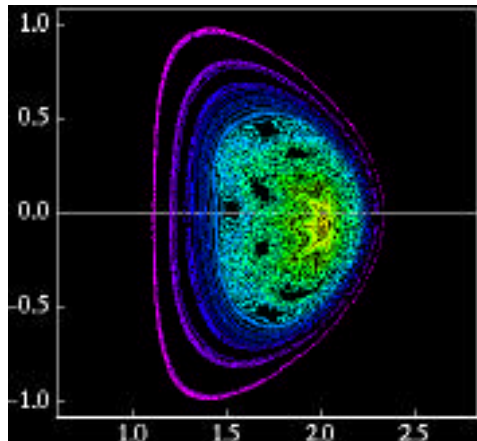
$t = 1.99 \times 10^{-4}$ sec



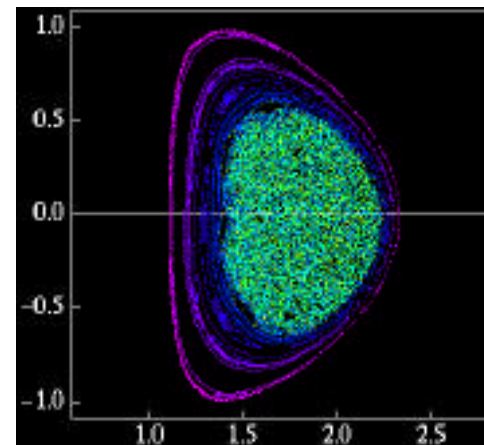
$t = 2.102 \times 10^{-4}$ sec



$t = 2.177 \times 10^{-4}$ sec

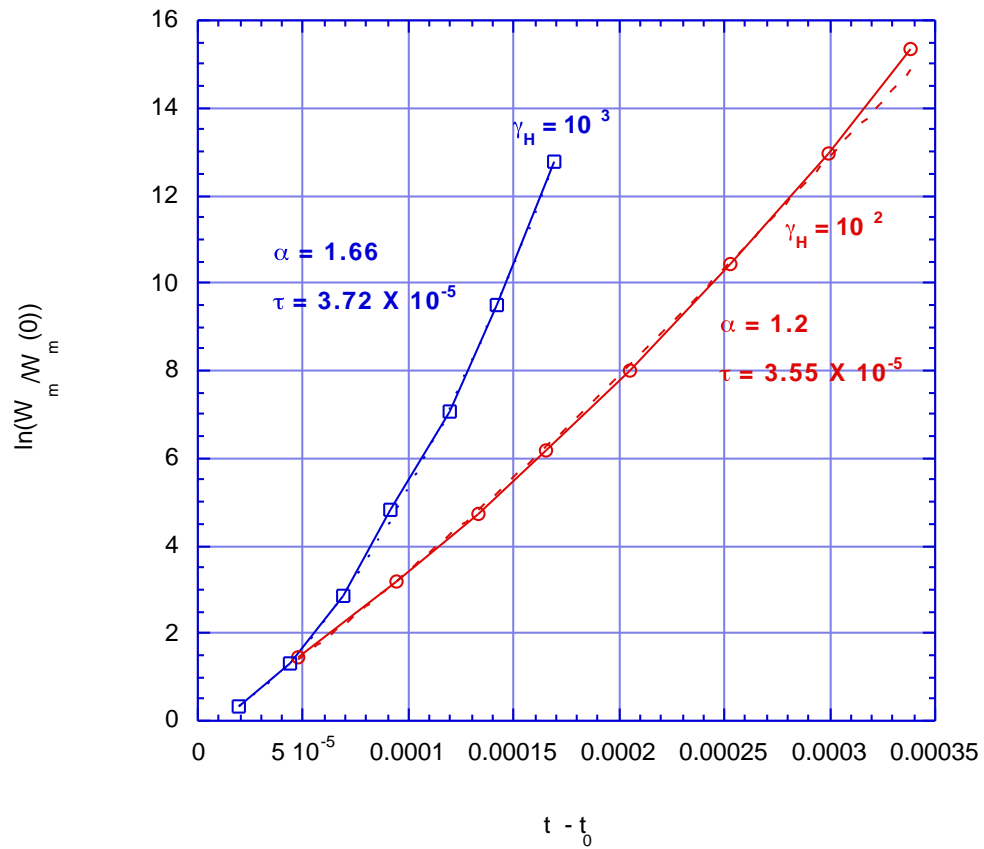


$t = 2.2 \times 10^{-4}$ sec



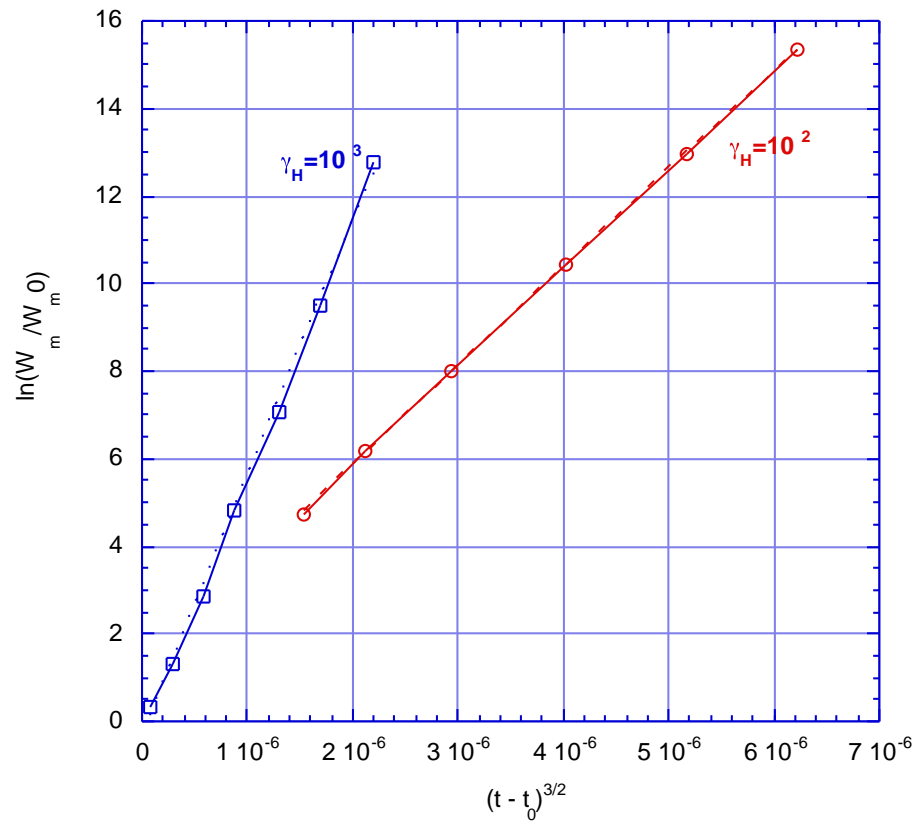
EFFECT OF HEATING RATE

$$\xi \sim \exp[(t - t_0)/\tau]^\alpha, \quad \alpha \text{ (theory)} = 3/2, \quad \tau \sim \gamma_{\text{MHD}}^{-2/3} \gamma_{\text{H}}^{-1/3}$$



ASSUME $\alpha = 3/2$ DEPENDENCE

$$\xi \sim \exp[(t - t_0)/\tau]^{3/2}, \quad \tau \sim \gamma_{\text{MHD}}^{-0.72} \gamma_{\text{H}}^{-0.28}$$



DISCUSSION

- Direct comparison with experiment difficult because of high β -limit due to presence of conducting wall
- Theory predicts $\xi = \xi_0 \exp \frac{t-t_0}{\tau}^{3/2}$, $\tau = \frac{3}{2} \hat{\gamma}_{MHD}^{-2/3} \gamma_H^{-1/3}$
- NIMROD finds $\xi = \xi_0 \exp \frac{t-t_0}{\tau}^{\alpha(\gamma_H)}$, $\tau \sim const.$
- When fit to $\alpha = 3/2$, NIMROD yields scaling $\tau \sim \gamma_{MHD}^{-0.72} \gamma_H^{-0.28}$
- Possible explanations
 - Initial conditions actually above critical β , then $\gamma_H \rightarrow 0$ would imply $\alpha \rightarrow 1$ (exponential growth)
 - Non-ideal effects are important
- Stochastic field lines in saturated state ==> degraded confinement
 - Self-consistent calculations with anisotropic heat flux underway